



Extreme Light Infrastructure: Attosecond Physics to Relativistic and Ultra-Relativistic Optics

*Roumanian Institute of Atomic
Physics*

20/11/2008

Gérard A. MOUROU

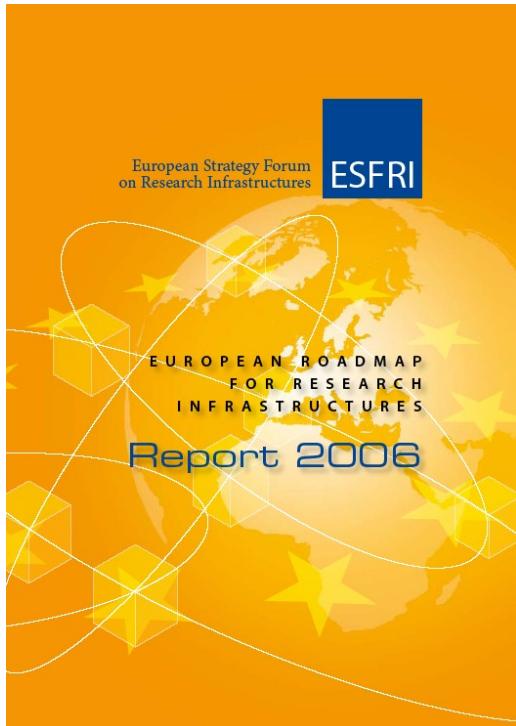
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Optics Horizon

This field does not seem to have natural limits, only horizon.

Two large Laser Infrastructures Have Been Selected to be on the ESFRI (European Strategic Forum on Research Infrastructures) Roadmap



- "a - HIPER, civilian laser fusion research (using the fast ignition scheme) and all applications of ultra high energy laser
- "b - ELI, reaching highest intensities (Exawatt) and applications

ELI has been the first Infrastructure launched by Brussels November 1st 2007



*Why should we build an Extreme
Light
Infrastructure?*

Science (1 July 2005)



“100 questions spanning the science...”

- 1) Is ours the only universe?
- 2) What drove cosmic inflation?
- 3) When and how did the first stars and galaxies form?
- 4) Where do ultrahigh-energy cosmic rays come from?
- 5) What powers quasars?
- 6) What is the nature of black holes?
- 7) Why is there more matter than antimatter?
- 8) Does the proton decay?
- 9) What is the nature of gravity?
- 10) Why is time different from other dimensions?
- 11) Are there smaller building blocks than quarks?
- 12) Are neutrinos their own antiparticles?
- 13) Is there a unified theory explaining all correlated electron systems?
- 14) **What is the most powerful laser researchers can build?** Theorists say an intense enough laser field would rip photons into electron-positron pairs, dousing the beam. But no one knows whether it's possible to reach that point.
- 15) Can researchers make a perfect optical lens?
- 16) Is it possible to create magnetic semiconductors that work at room temperature?

Contents

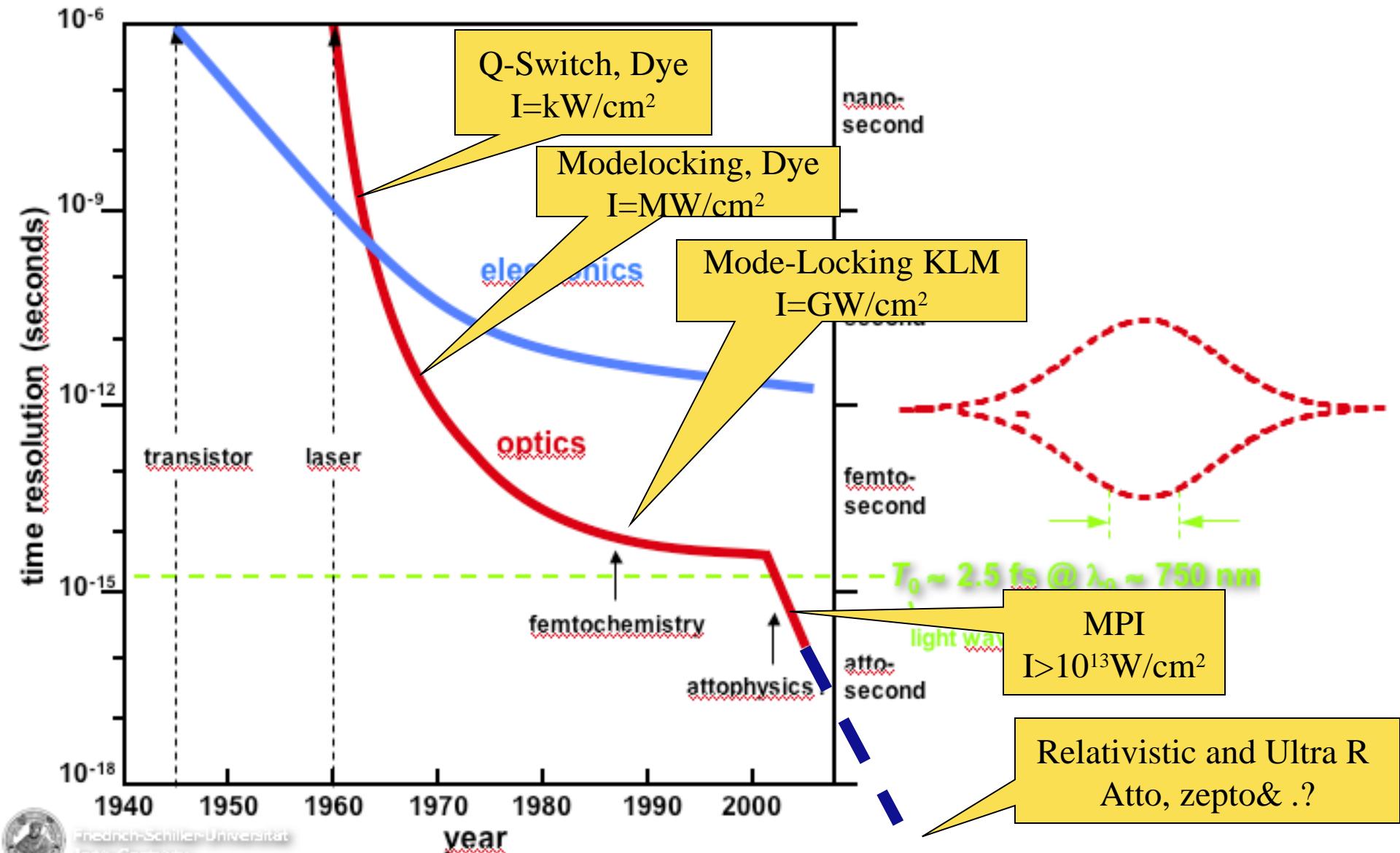
- The Peak Power-Pulse Duration conjecture
- Relativistic Optics: A parallel with Bound Electron Nonlinear Optics
- Relativistic Rectification(wake-field) the key to High energy electron beam, proton beam, x-ray and γ -beams
- Source of attosecond photon and electron pulses
- Generation of Coherent x-ray by Coherent Thomson scattering
- A route to the critical (Schwinger)field
- Few examples of applications to High energy Physics
General Relativity, Hawking Radiation, Extradimensions.

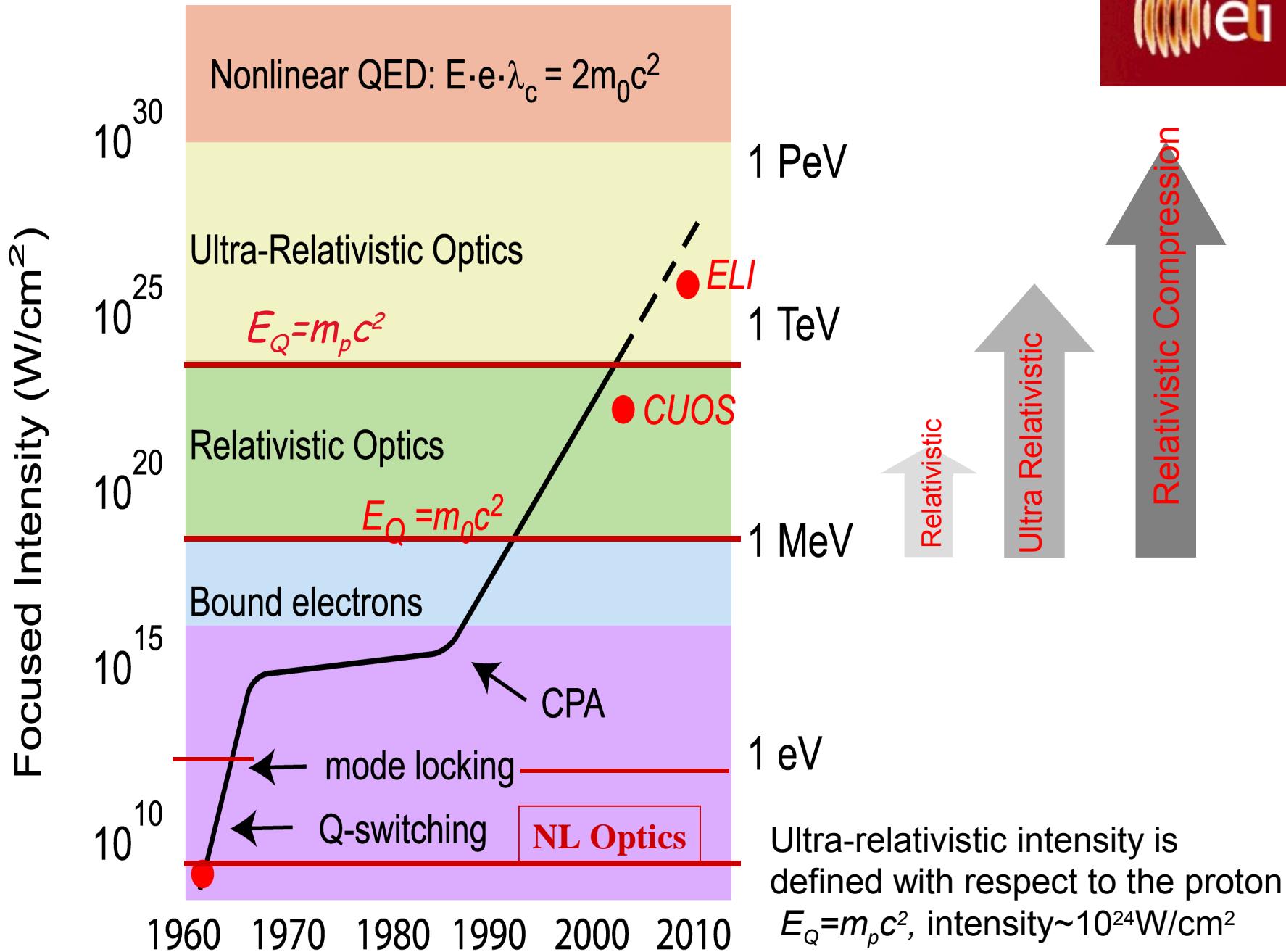
ELI(Extreme Light Infrastructure) An Exawatt laser Infrastructure on the Large Infrastructure Road Map of Europe

Peak Power -Pulse Duration Conjecture

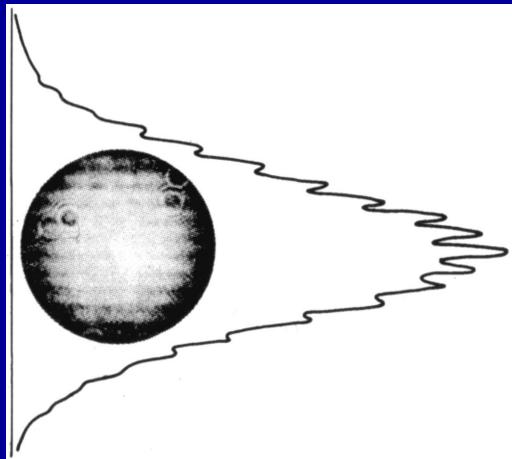
- 1) To get high peak power you must decrease the pulse duration.
- 2) To get short pulses you must increase the intensity

Laser Pulse Duration vs. Intensity

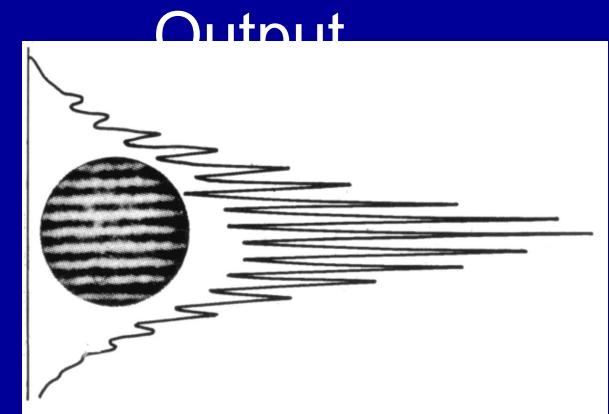




Small-Scale Self-Focusing



Amplifier
 $n = n_0 + n_2 I$



Instabilities grow with a maximum growth rate:

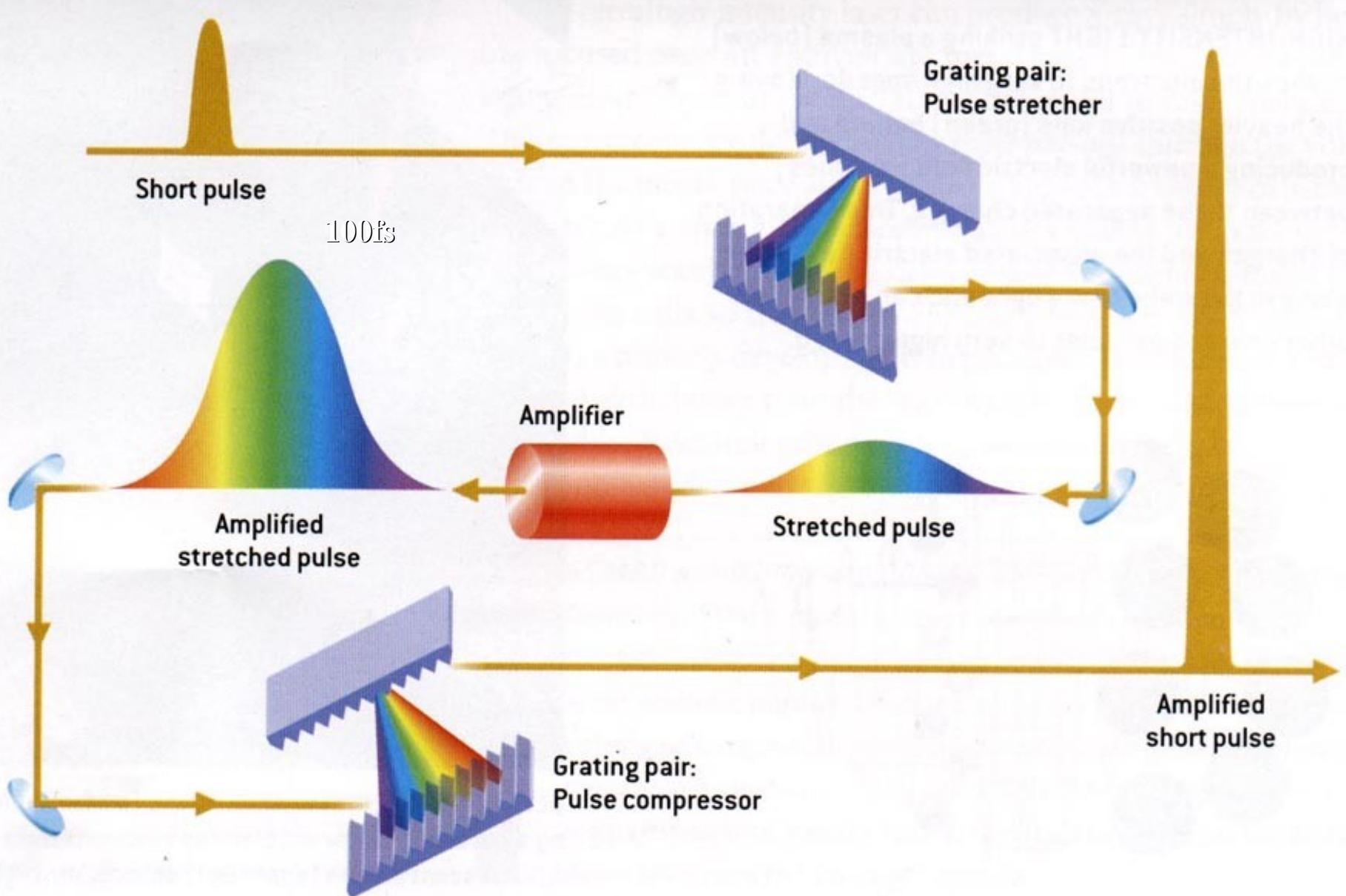


B-integral < 3 for good beam quality:

$$B = \frac{2p}{l} \int_0^L n_2 I(z) dz$$

Chirped Pulse Amplification

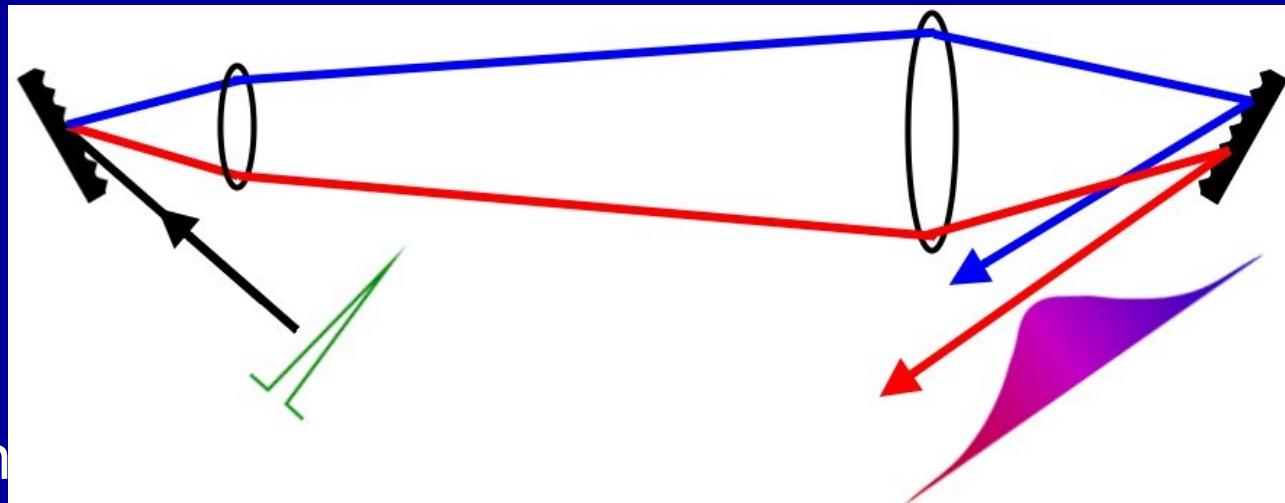
D. Strickland and G. Mourou 1985



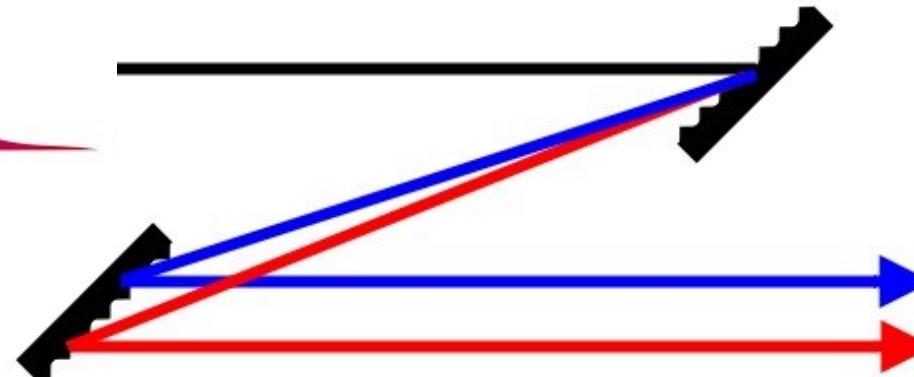
Matched Stretcher-Compressor

1000 Times Expansion/Compression of Optical Pulses for Chirped Pulse Amplification

M. Pessot, P. Maine, and G. Mourou, *Optics Commun.* **62**, 419-421 (June 1987)



Stretch



Compressor

Bound Electron Nonlinear Optics

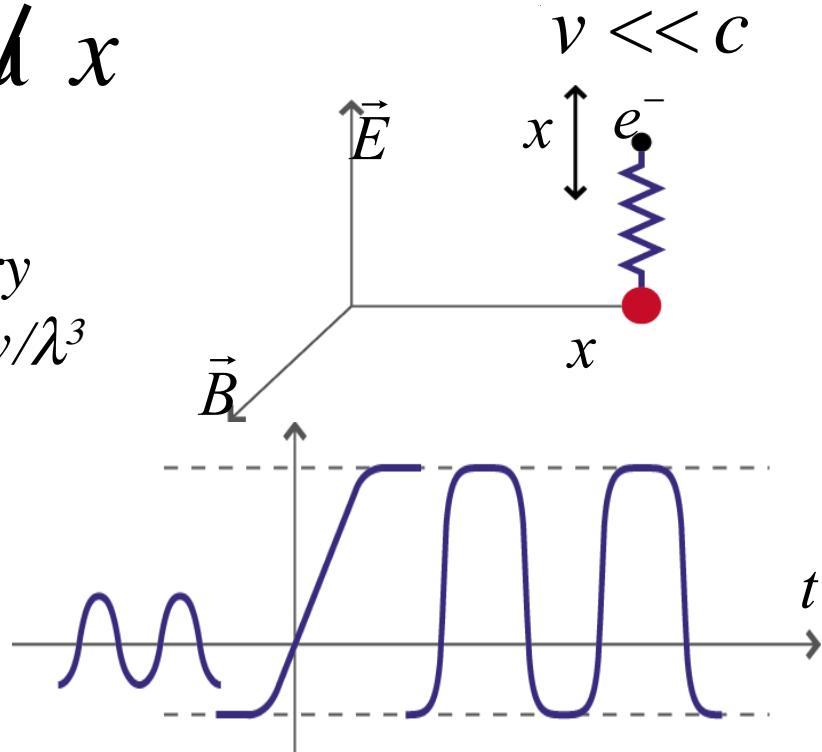


PETER B. HIRSCH



$$\vec{F} = q\vec{E} \quad F \not\propto x$$

*The field necessary
corresponds to hv/λ^3*



"*Harmonics*

"*Optical Rectification*

"*Self-focusing*

Relativistic Optics

Relativistic Optics

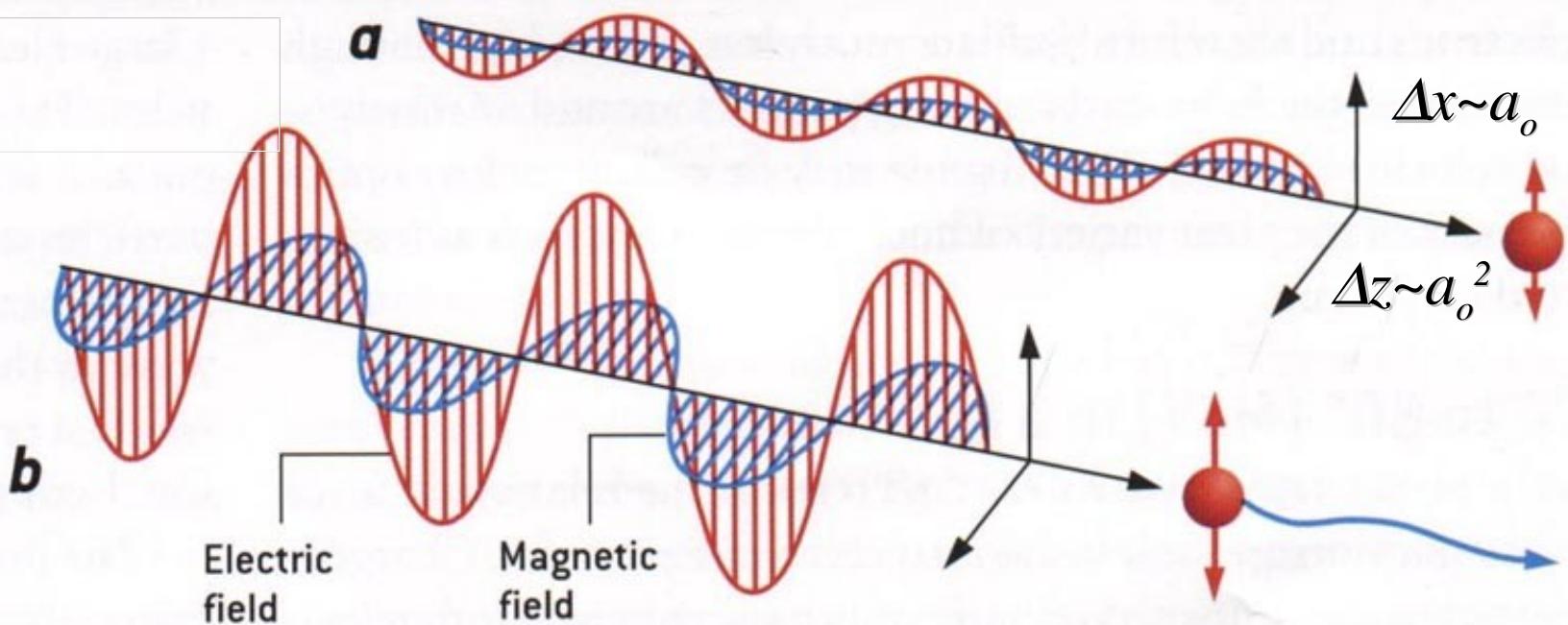


a) Classical optics $v \ll c$,

$$a_0 \ll 1, a_0 \gg a_0^2$$

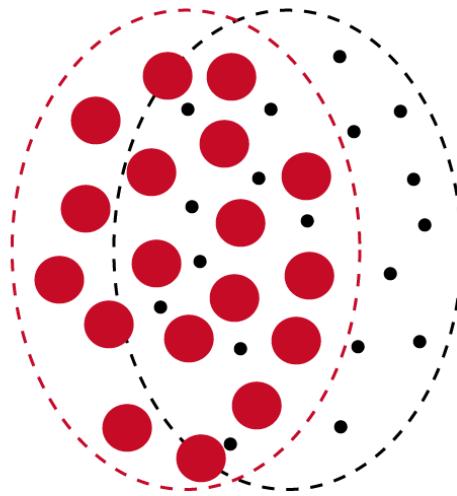
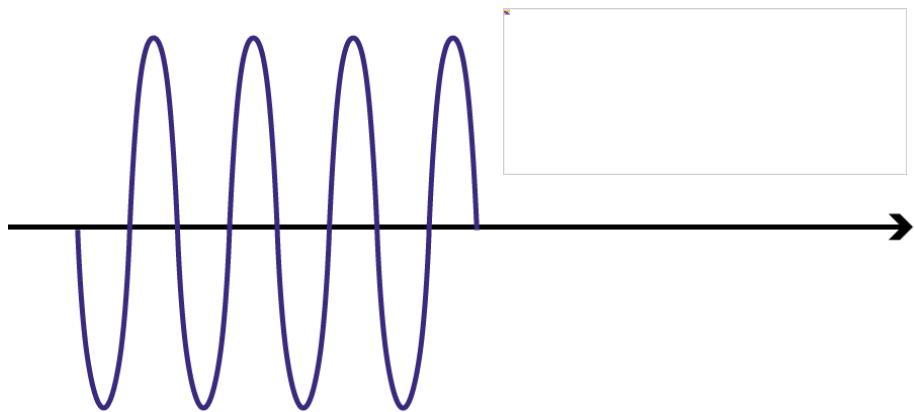
b) Relativistic optics $v \sim c$

$$a_0 \gg 1, a_0 \ll a_0^2$$



Relativistic Rectification

(*Wake-Field Tajima, Dawson*) \vec{E}_s 



- 1)  pushes the electrons.
- 2) The charge separation generates an electrostatic longitudinal field. (Tajima and Dawson: Wake Fields or Snow Plough)

- 3) The electrostatic field


Relativistic Rectification

-Ultrahigh Intensity Laser is associated with Extremely large E field.

$$E_L^2 = Z_0 * I_L$$

← →
Medium Impedance Laser Intensity

$$I_L = 10^{18} W/cm^2$$

$$E_L = 2 \quad TV/m$$

$$I_L = 10^{23} W/cm^2$$

$$E_L = .6 PV/m \quad (0.6 \cdot 10^{15} V/m)$$

Laser Acceleration:

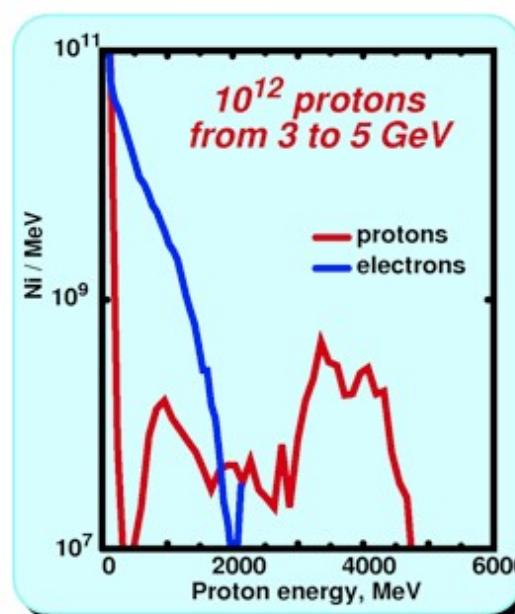
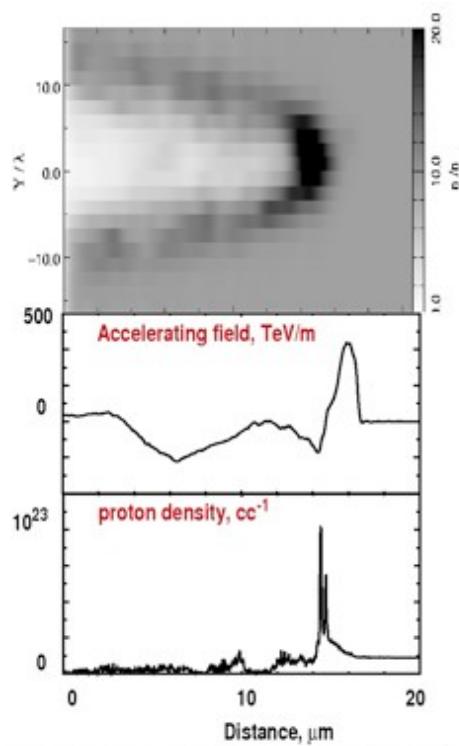
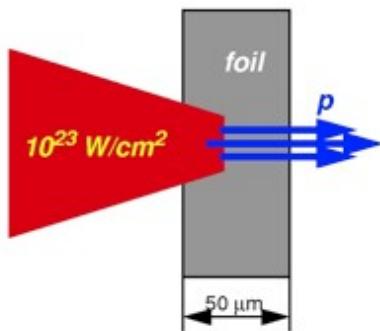
At 10^{23}W/cm^2 , $E= 0.6 \text{PV/m}$, it is SLAC (50GeV, 3km long) on $10\mu\text{m}$ The size of the Fermi accelerator will only be one meter (PeV accelerator that will go around the globe, based on conventional technology).

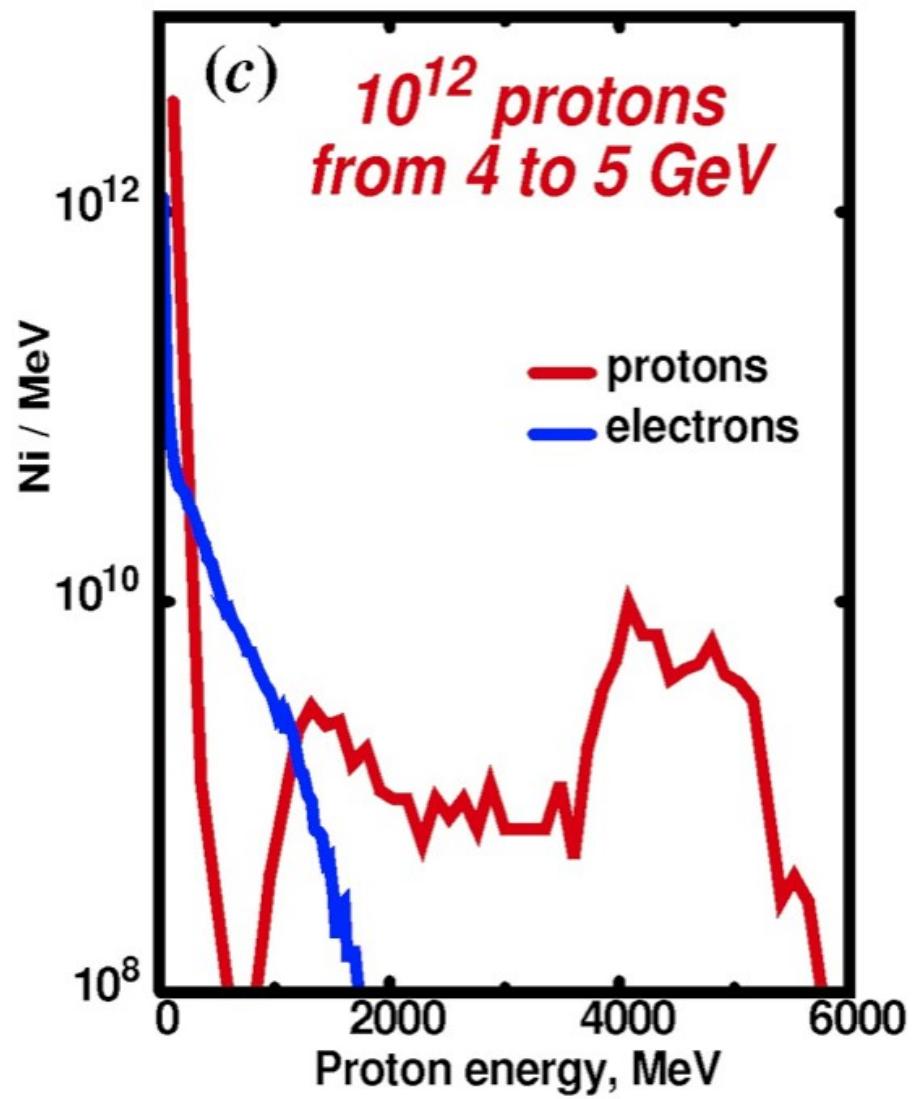
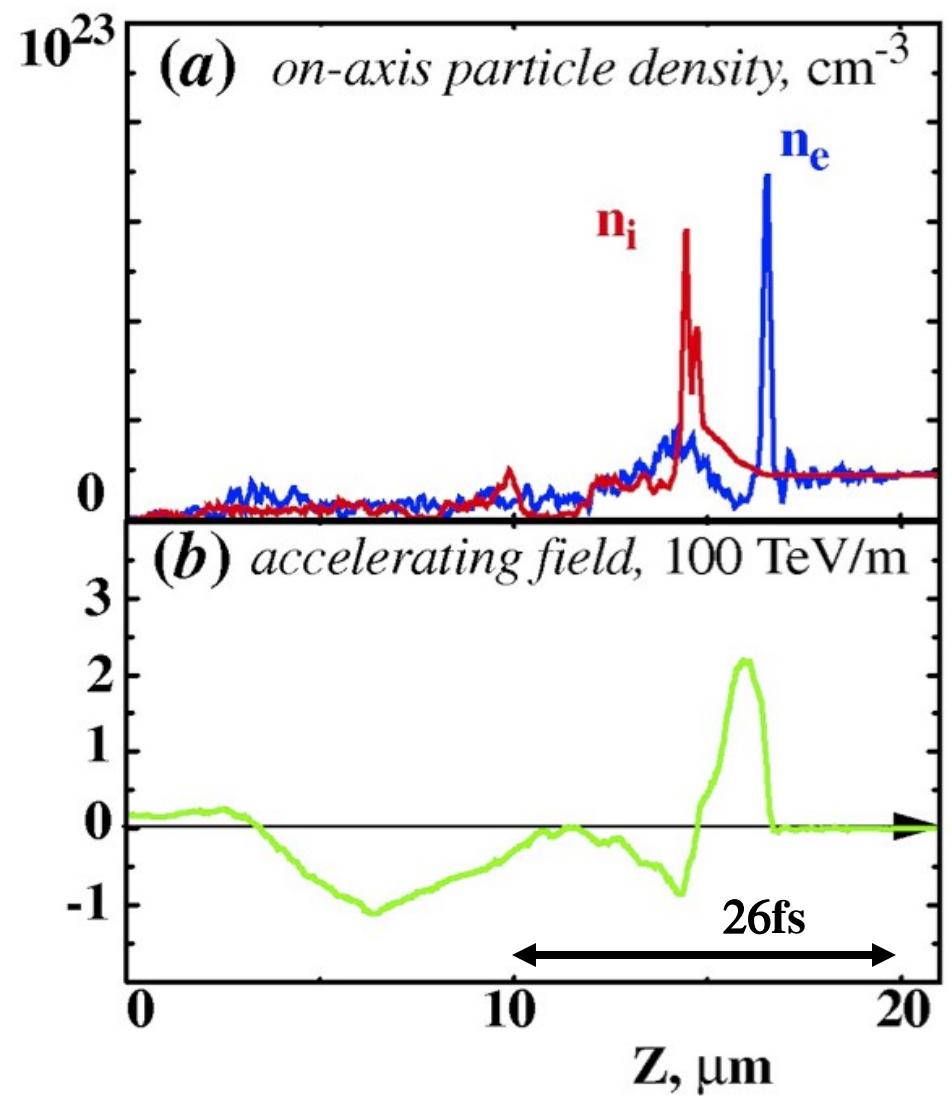
Relativistic Microelectronics



5 GeV proton bunch at solid state density

3d PIC simulations, A.Pukhov, Theorie, MPQ,





The Dream Beam

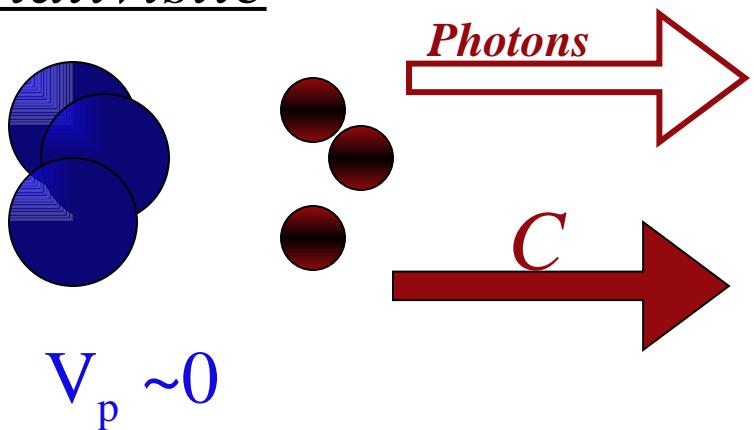


J. Faure et al., C. Geddes et al., S. Mangels et al. ,
in *Nature* 30 septembre 2004

Relativistic Protons

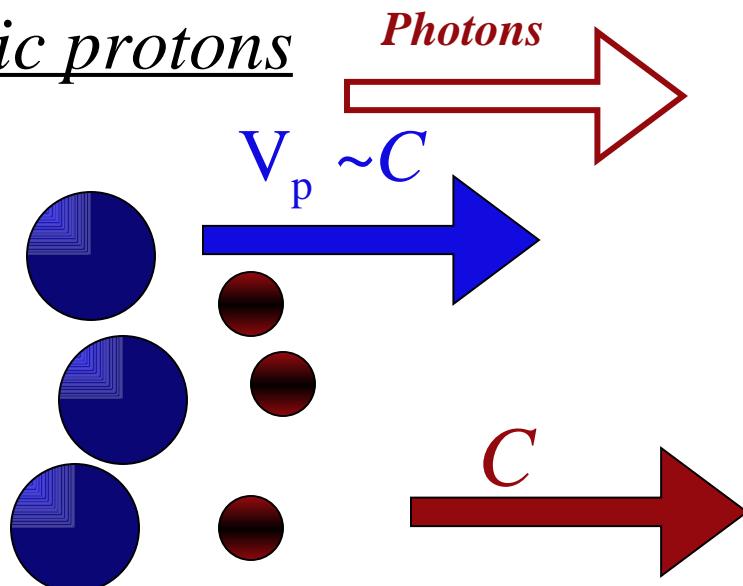


Non relativistic



$$E_p \sim I^{1/2}$$

Relativistic protons

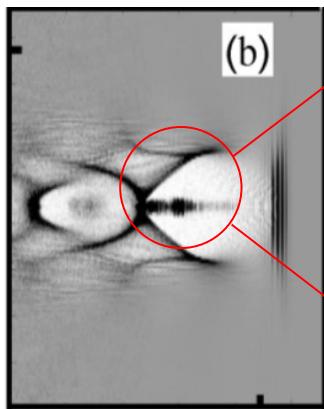
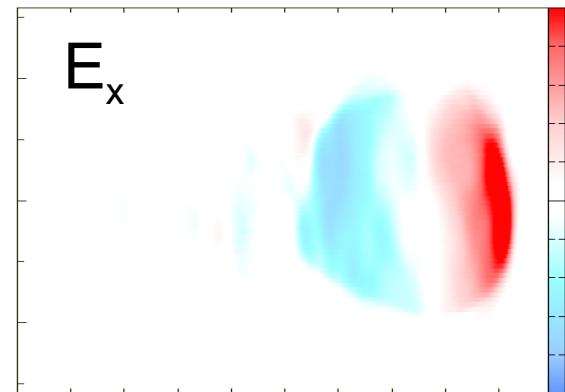
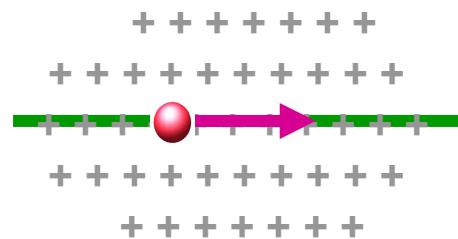


$$E_p \sim I$$

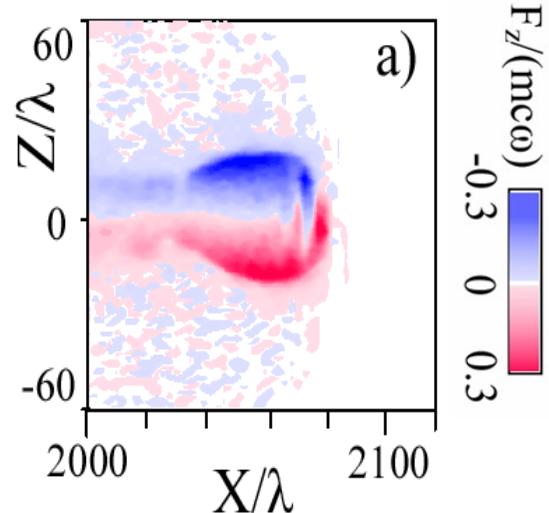
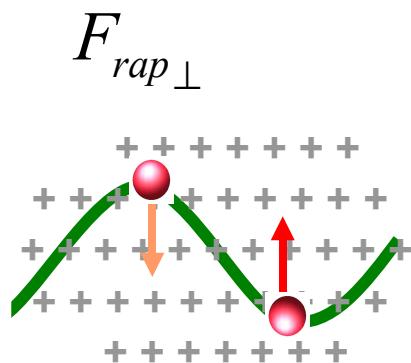
The structure of the ion cavity



Longitudinal acceleration



Transverse oscillation: Betatron oscillation

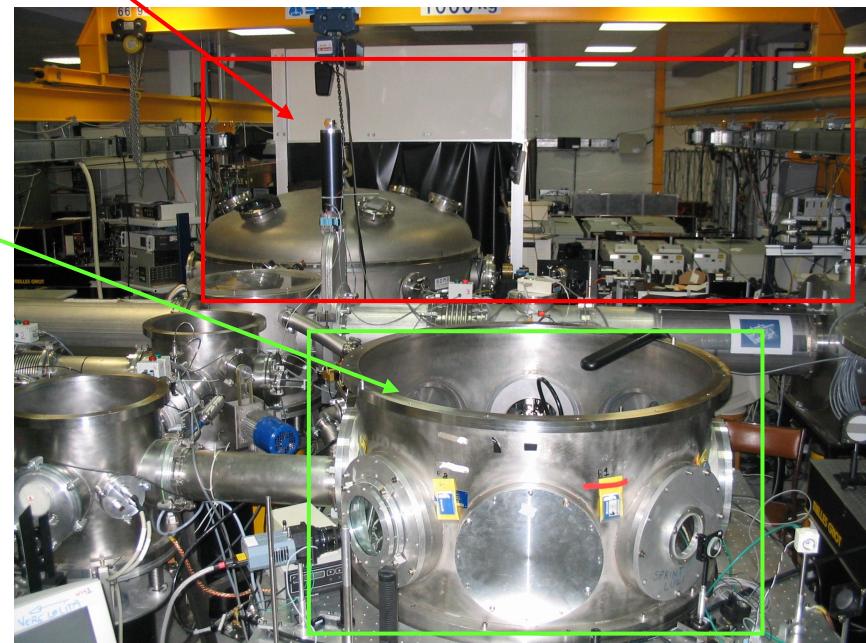
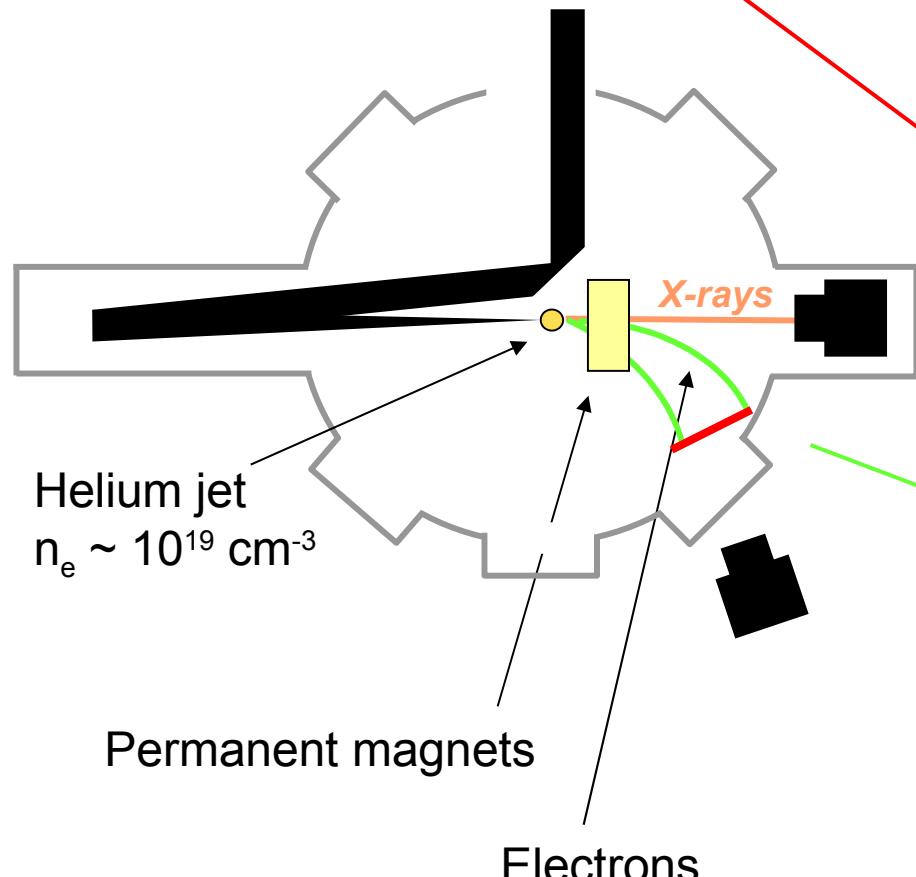


Experiment Setup

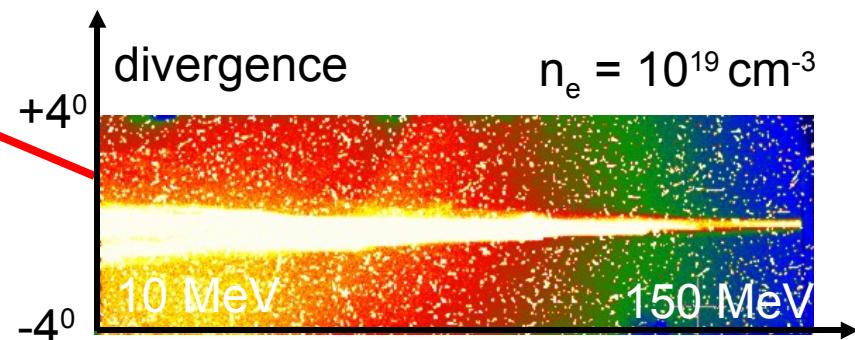
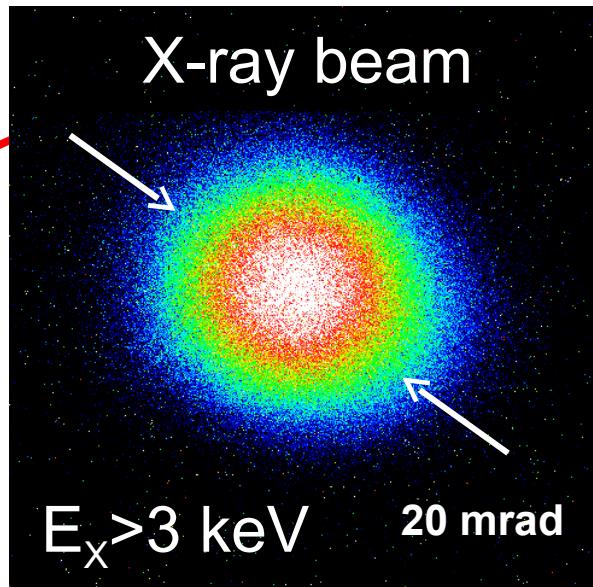
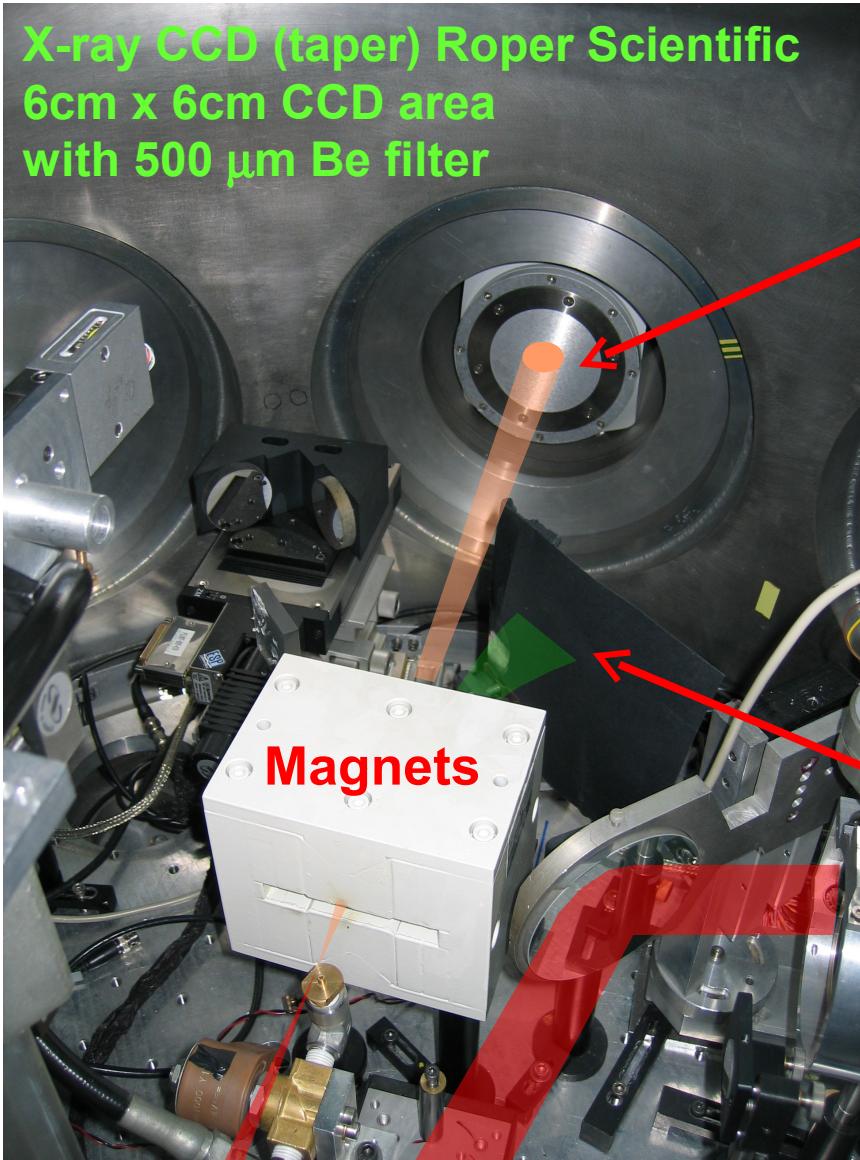
A. Rousse et al.



Laser 1.5 J/30 fs (Salle Jaune)



Simultaneous measurements of X-ray and Electron Beams



Electron beam
K. Ta Phuoc et al

Attosecond Generation (photon)

Attosecond by Bond electron Nonlinear Optics

- P. Corkum, M. Ivanov and Burnett Sub.femtosecond Pulses Opt. Lett. 19, 1870 (1994)
- M. Hentschel *et al.* Nature 414, 509 (2001)
- A. Baltuslka *et al.* Nature 421, 6111 (2001)

The technique relies on High harmonic Generation and not very efficient.

It is limited to nJ and not scalable to high energy.

HHG and Subfemtosecond Pulses from Surfaces of Overdense Plasmas

**S.V. Bulanov, Naumova N M and Pegoraro F, Phys. Plasmas
1 745(1994)**

D. Von der linde et al Phys. Rev. A52 R 25, 1995

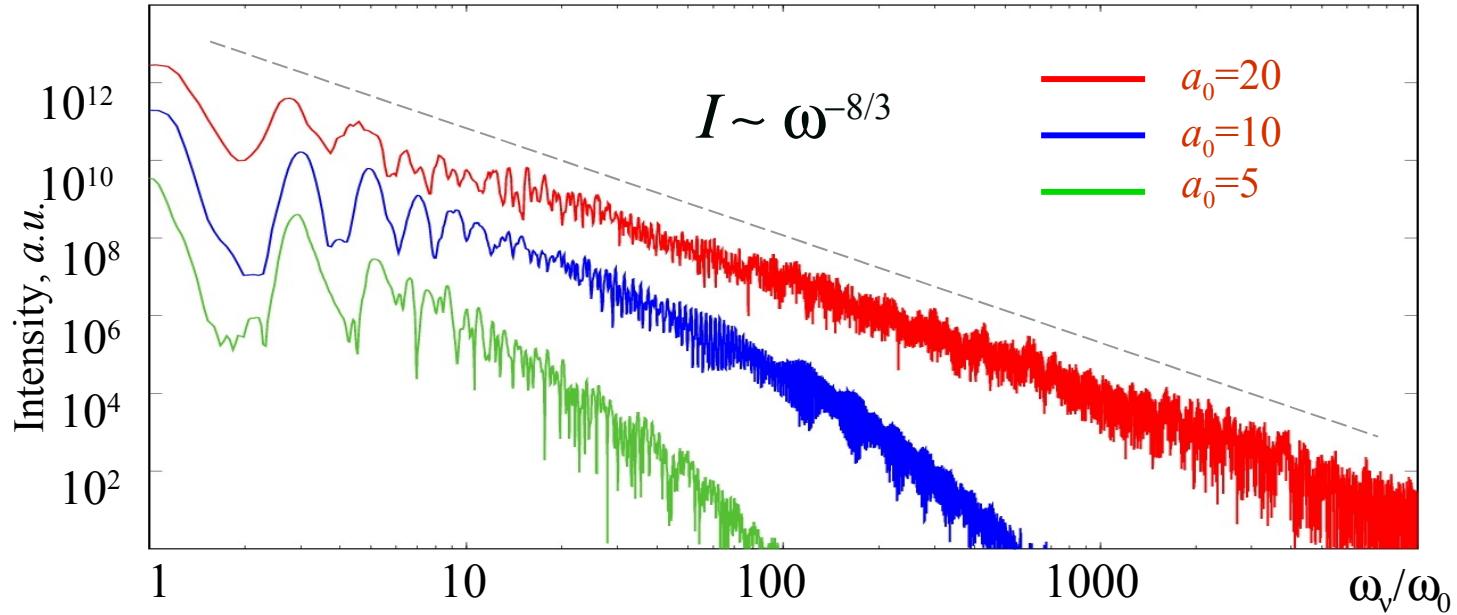
L. Plaja et al. JOSA B, 15, 1904 (1998)

S. Gordienko et al PRL 93, 115002 (2004)

N.M. Naumova et.al., PRL 92, 063902 (2004)

Tsakiris, G., *et al.*, New Journal of Physics, 8, 19 (2006)

Reflected radiation spectra: the slow power-law decay 1D simulation



Gordienko, et al., Phys. Rev. Lett. 2004

The Gaussian laser pulse $a=a_0\exp[-(t/\tau)^2]\cos\omega_0 t$ is incident onto an overdense plasma layer with $n=30n_c$.

The color lines correspond to laser amplitudes $a_0=5,10,20$.

The broken line marks the analytical scaling $I \sim \omega^{-8/3}$.

Possibility to produce zeptosecond pulses!!!

VULCAN Experiment: Harmonics down to “Water Window”



B. DROMEY, M. ZEPF, et al., **NATURE Physics**, Vol. **2**, p. 456 (2006).

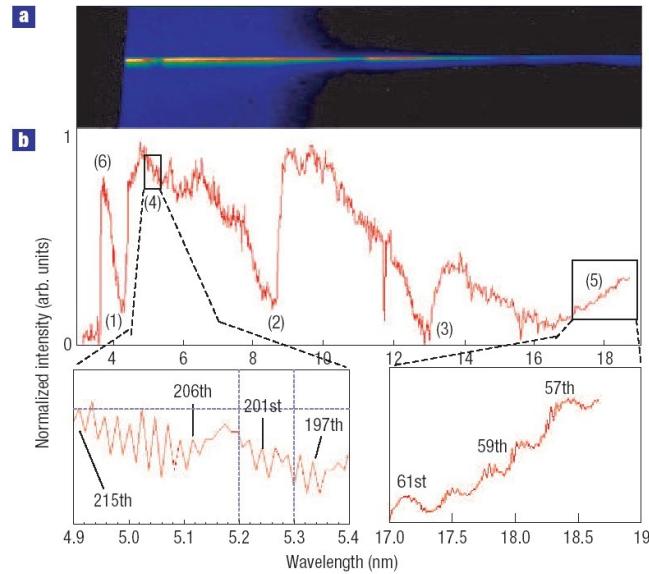


Figure 3 Unprocessed high harmonic spectrum recorded with the extreme-ultraviolet spectrometer. **a**, Raw CCD image obtained with the double PM setup ($E = 70$ J on target, false colours). **b**, A lineout of **a**. Spectral features: (1) first-order carbon K-edge (4.36 nm), (2) second-order carbon K-edge, (3) third-order carbon K-edge, (4) region of resolved harmonics around 200th order

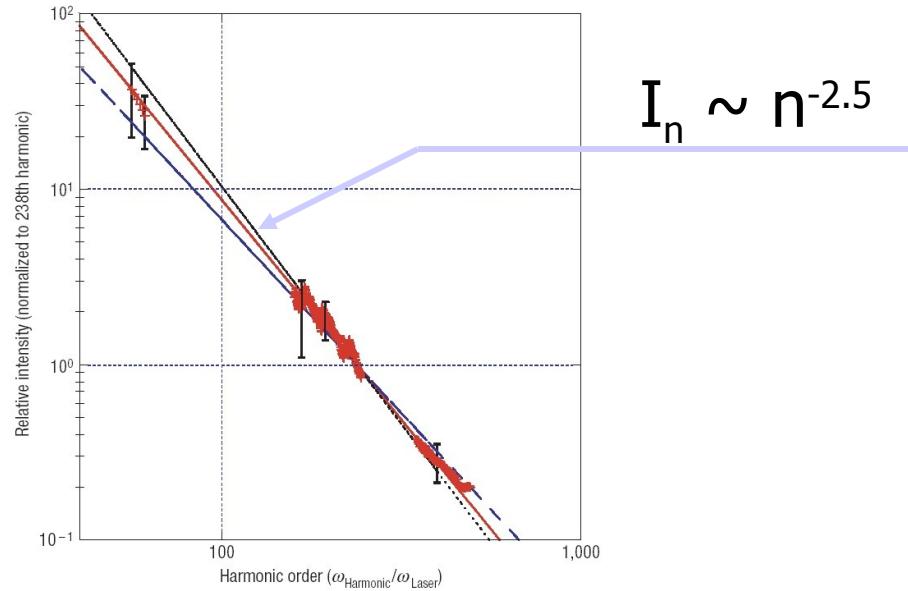


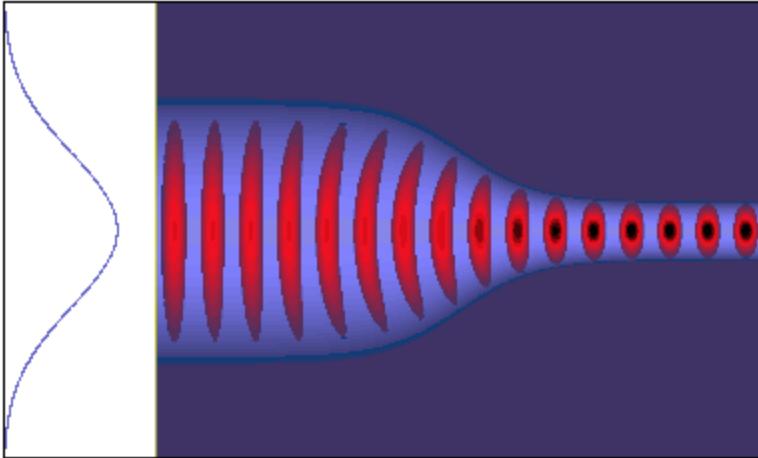
Figure 4 Relative intensity of harmonics normalized to the 238th harmonic (at the carbon K-edge). The lines are fits to the data with the exponent p as a fitting parameter such that $I(n)/I(238) = n^{-p}/238^{-p}$. The best fit (red line) corresponds to a value of $p = 2.5$ confirming harmonic production in the relativistic limit. The error

Relativistic Self-focusing:

$$\varepsilon = 1 - \frac{w_p^2}{g_0 w^2}$$

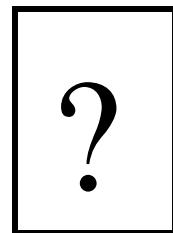
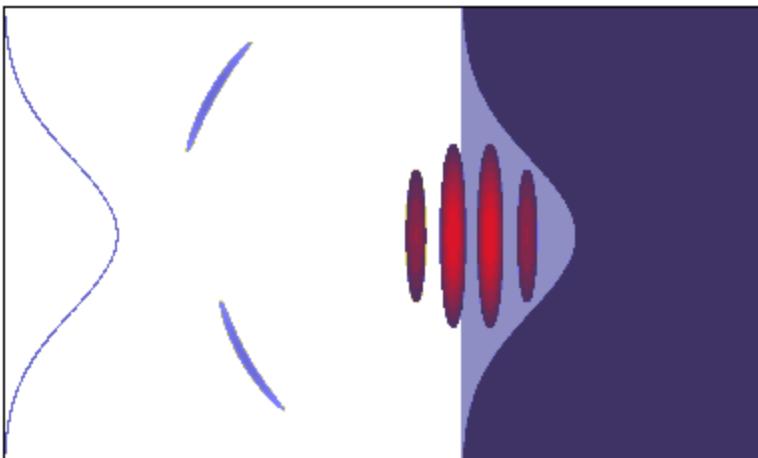
where $g_0 = \sqrt{1 + |a_0|^2}$

(a)
Refraction

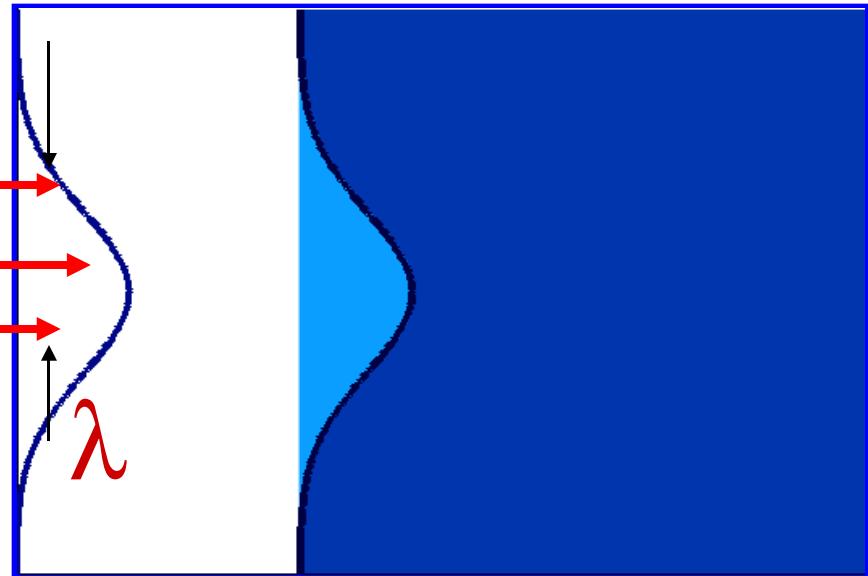


A.G.Litvak (1969), C.Max,
J.Arons, A.B.Langdon
(1974)

(b)
Reflection



Single Mode Relativistic optics in Reflection

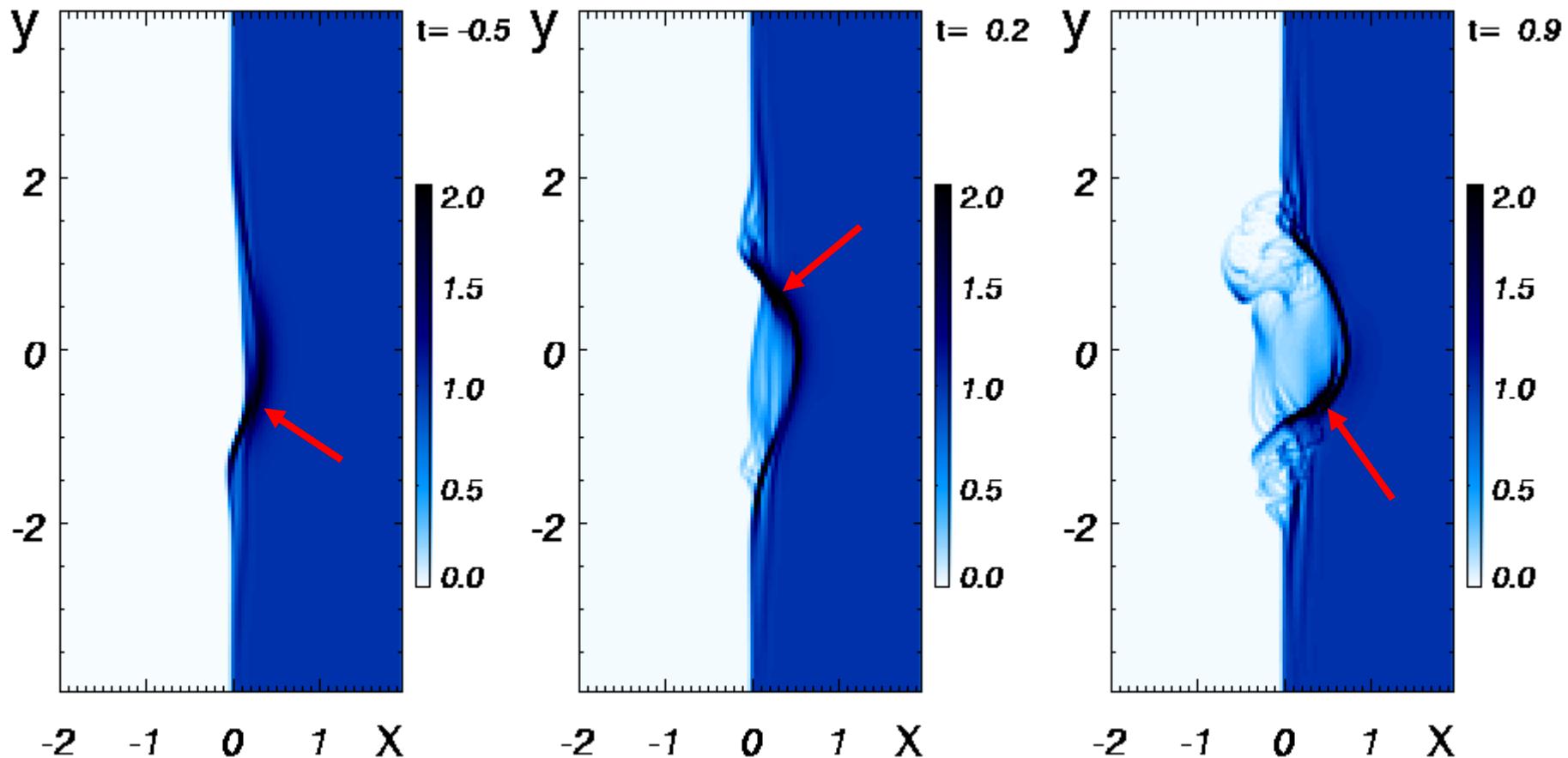


Under the action of the light pressure the critical surface will be pushed (curved) at relativistic speed at twice the laser frequency (ω).

If the laser is focused on 1λ , it will act as a perfect single mode mirror, leading to well behaved reflection and deflection.

The restoration force is a function of the plasma density

Moving plasma profile deflecting the isolated attosecond pulses at the instants of their generation

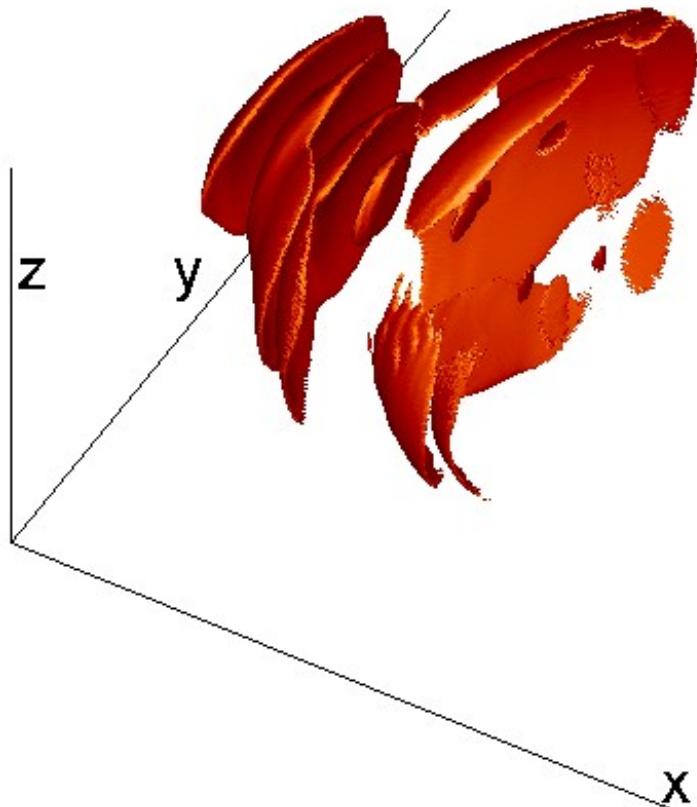


Relativistic electrons create the Doppler compression

N. M. Naumova, J. A. Nees, I. V. Sokolov, B. Hou, and G. A. Mourou, Relativistic generation of isolated attosecond pulses in a λ^3 focal volume, *Phys. Rev. Lett.* **92**, 063902-1 (2004).

3-D PIC simulation

*Electromagnetic
energy density*

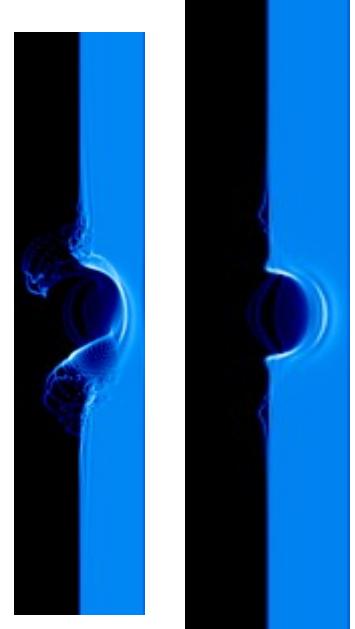


attosecond pulse

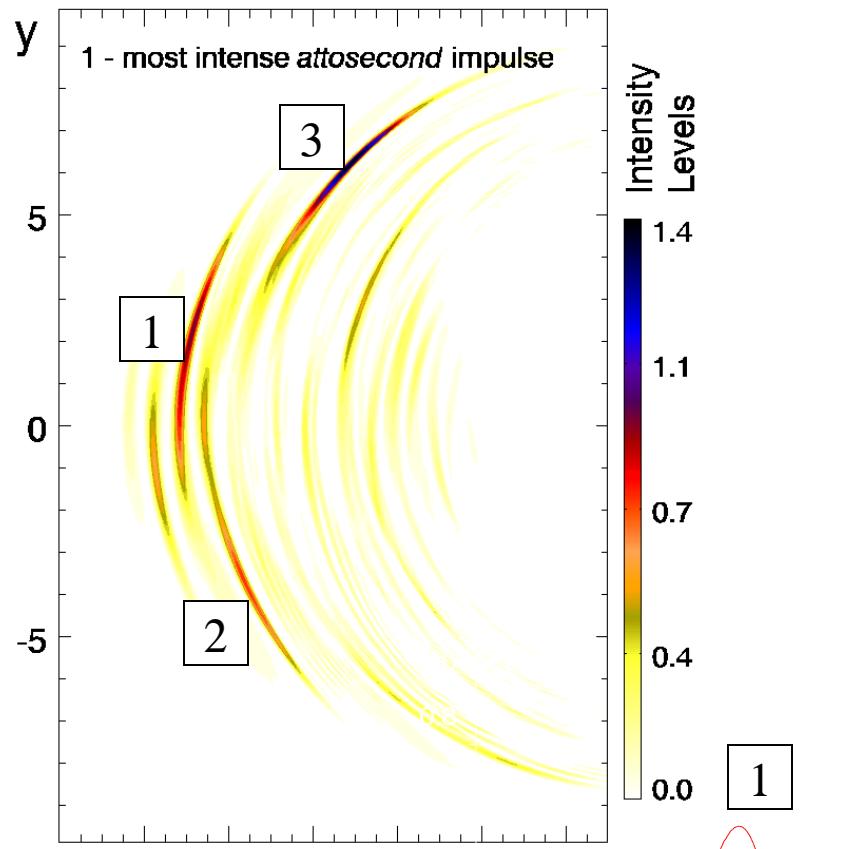


*Electron
density*

P-plane S-plane



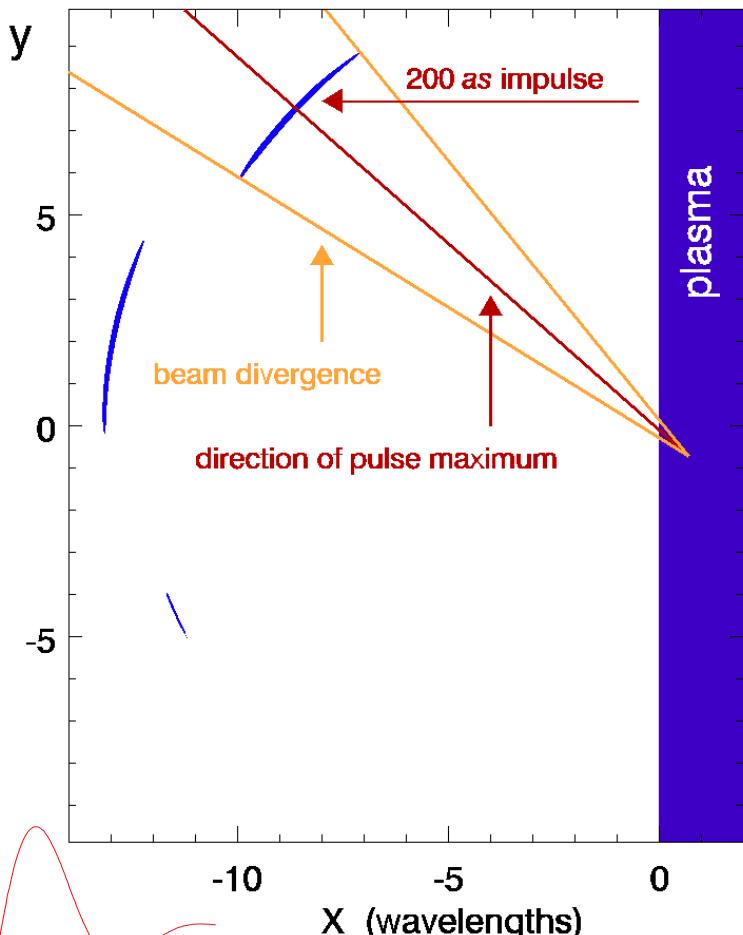
2-D PIC simulation



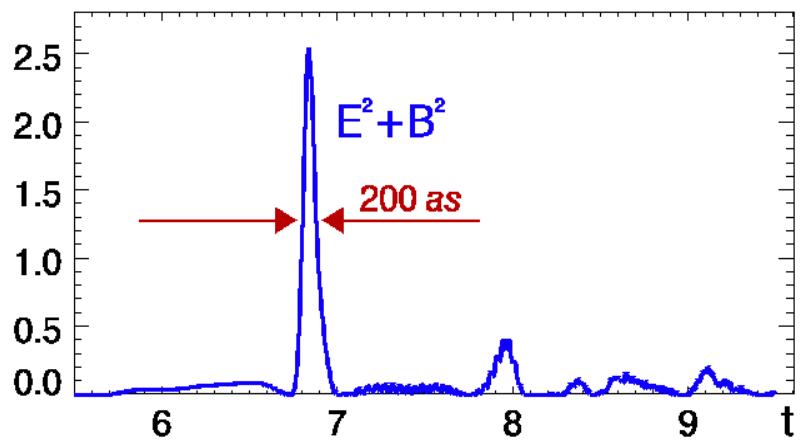
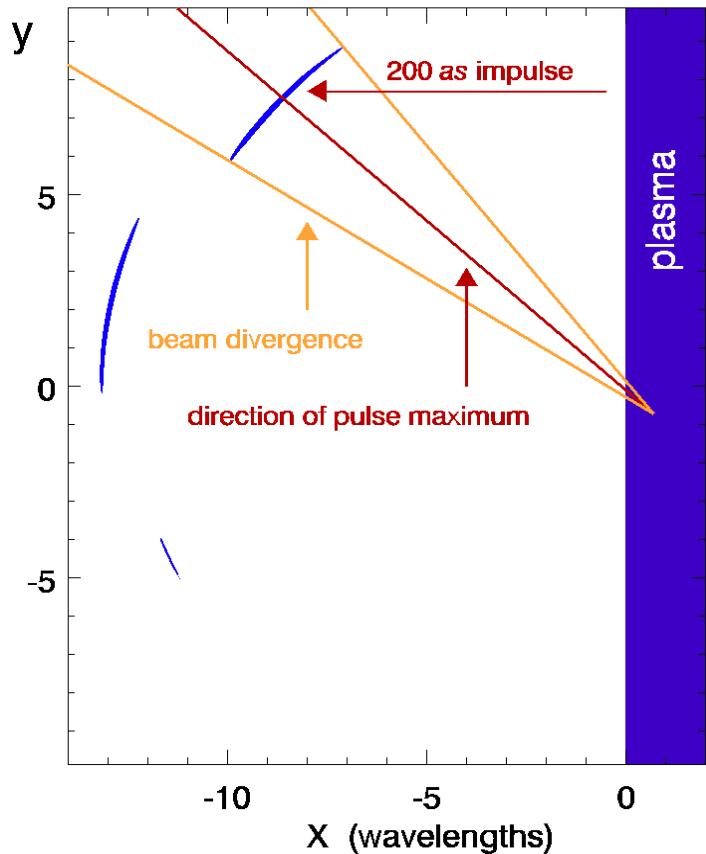
3

1

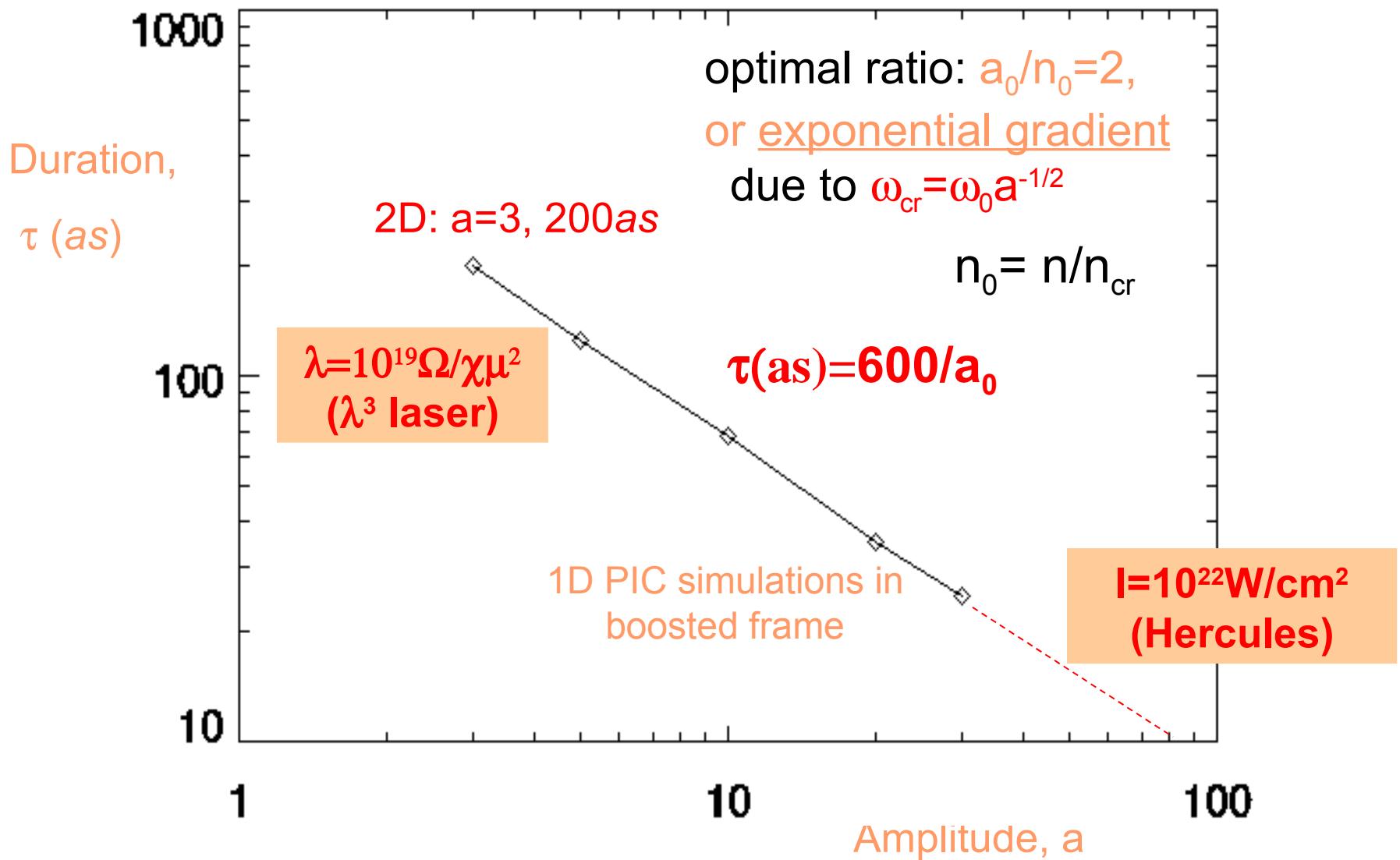
2



2-D PIC simulation

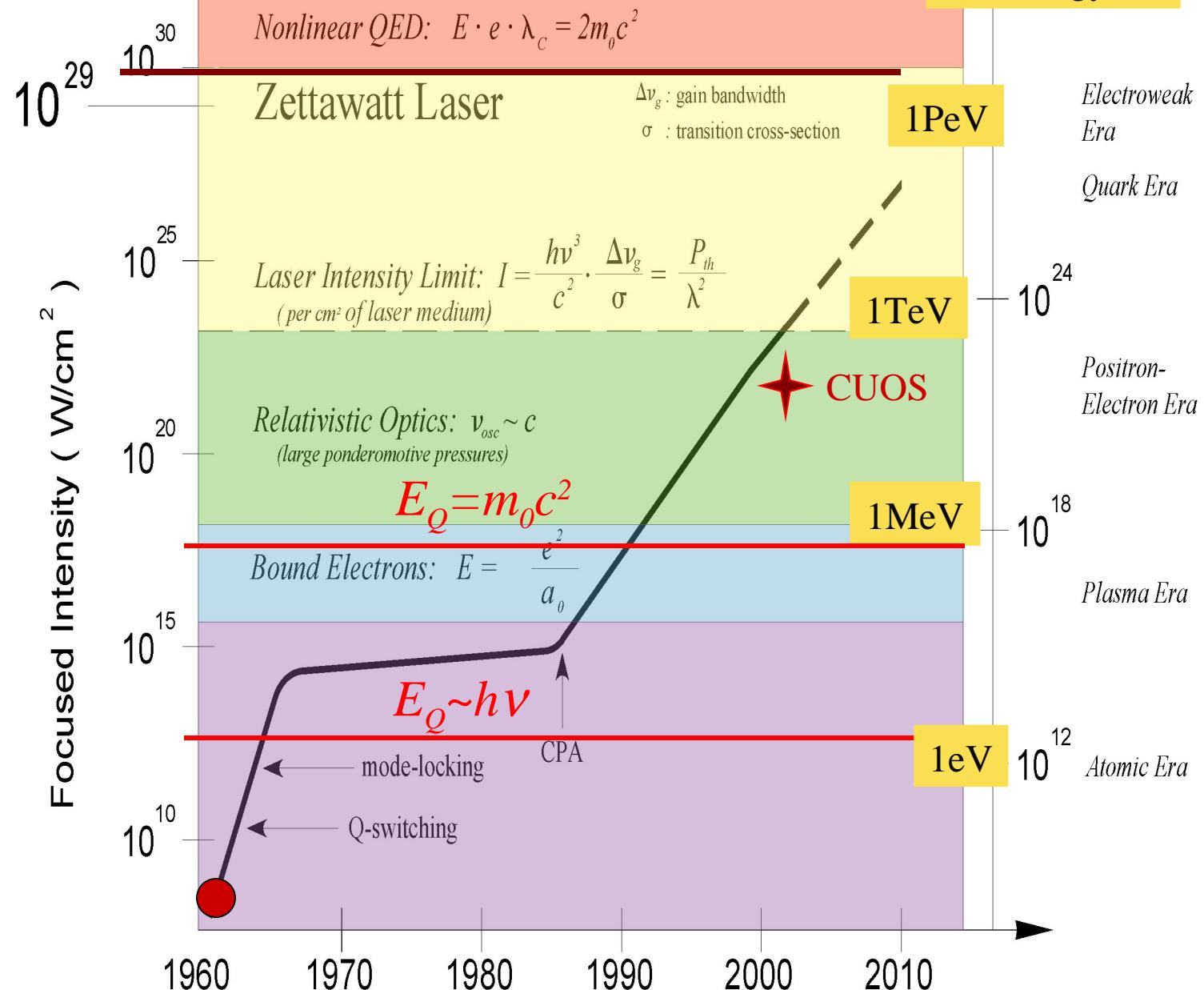


Scalable Isolated Attosecond Pulses





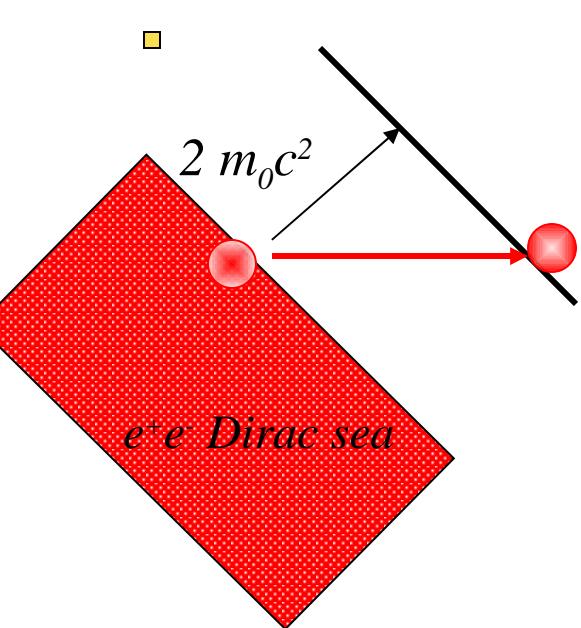
Electron Characteristic Energy



Laser-Induced Nonlinear QED

E. Brezin and C. Itzykson Phys. Rev.D2,1191(1970)

Vacuum can be considered like a dielectric



Schwinger Field

$$E_s = 1.3 \cdot 10^{16} \text{ V/cm}$$



Vacuum Tunneling

$$I_s = 10^{30} \text{ W/cm}^2$$



Towards the Critical Field

For $I=10^{22} \text{W/cm}^2$ $a_0^2 = 10^4$

The pulse duration $\tau = 600 / a_0 \sim 6 \text{as}$

The wavelength $\sim \lambda/1000$

The Focal volume decreases $\sim 10^{-8}$

The Efficiency $\sim 10\%$

Intensity $I=10^{22} \text{W/cm}^2 \longrightarrow I=10^{28} \text{W/cm}^2$

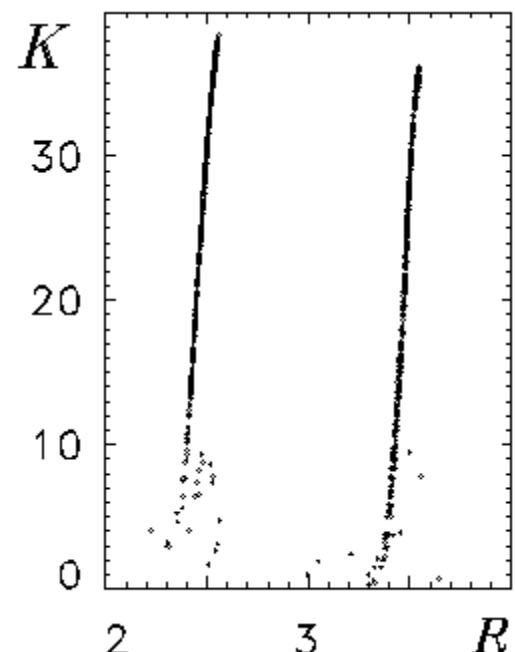
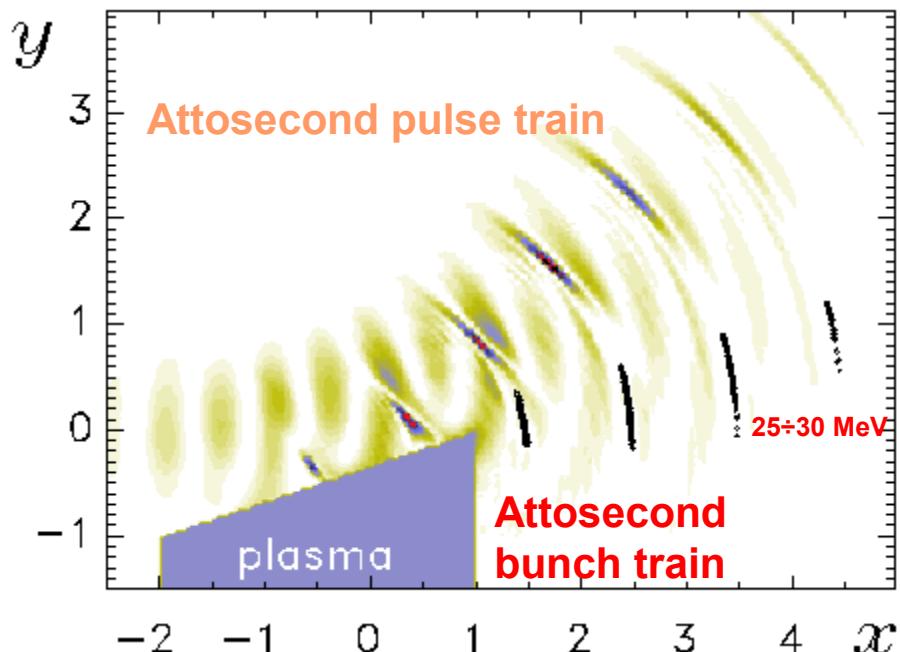
Isolated Attosecond Pulse Generation by Relativistic Compression and Deflection in the λ^3

- Optimum use of the energy
- Provides pulse isolation
- Predicted to be efficient (10^{-1} - 10^{-2})
- Single mode and short pulse suppress the instabilities, that is very good beam quality expected.
- It relies on relativistic plasma therefore scalable to any pulse energy (kJ)
- The higher the intensity the shorter the pulse

Attosecond Generation (electron)

Attosecond Electron Bunches

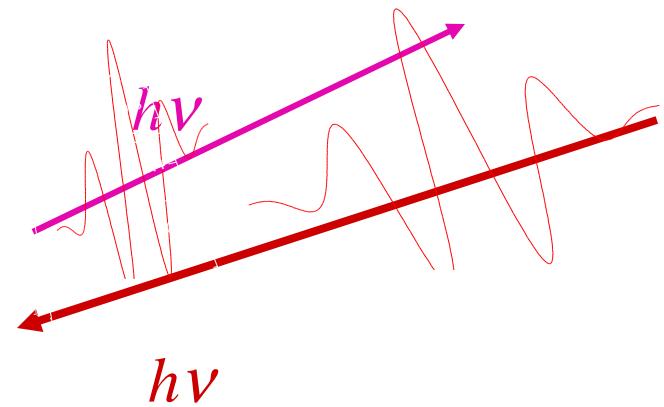
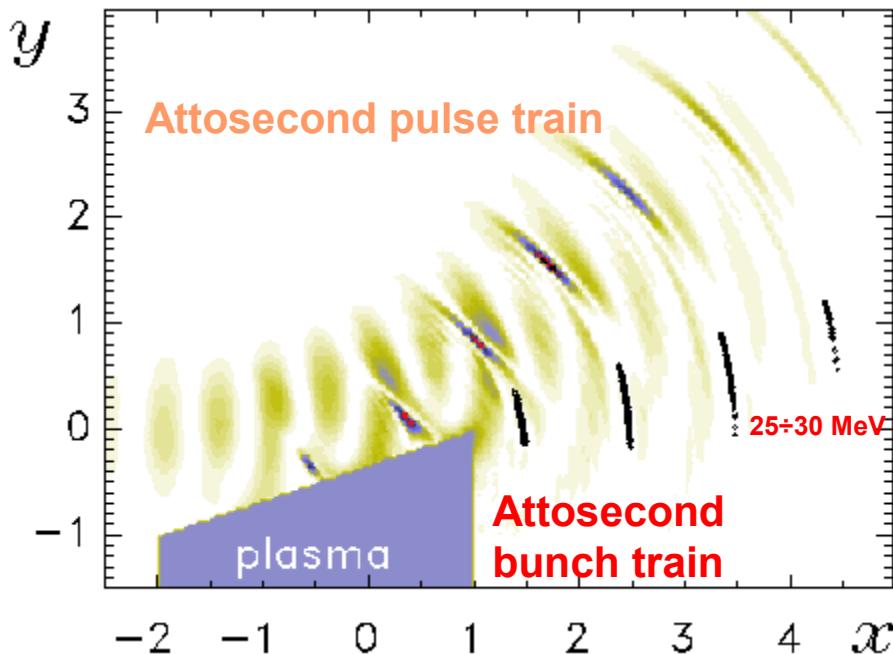
$$a_0=10, \tau=15\text{fs}, f/1, n_0=25n_{\text{cr}}$$



N. Naumova, I. Sokolov, J. Nees, A. Maksimchuk, V. Yanovsky, and G. Mourou,
 Attosecond Electron Bunches, *Phys. Rev. Lett.* **93**, 195003 (2004).

Coherent Thomson Scattering

$$a_0=10, \tau=15\text{fs}, f/1, n_0=25n_{\text{cr}}$$



N. Naumova, I. Sokolov, J. Nees, A. Maksimchuk, V. Yanovsky, and G. Mourou,
Attosecond Electron Bunches, *Phys. Rev. Lett.* **93**, 195003 (2004).

Electron bunches of ~100 as duration would produce backward

Coherent Thomson scattering efficiency

- Cross-section for the backward Thomson scattering:

$$\sim N + N(N-1) \exp(-2(k'd')^2)$$

depends on the factor in the exponent: $k'd' = kd(1 + V/c)^2 \gamma^2$.

- The resulting backward Thomson cross-section

$$\sigma_T N^2 \exp(-8(kd)^2 \gamma^4) \sim 10^{-4} \exp(-8(kd)^2 \gamma^4) \text{ cm}^2$$

is far above the channel cross-section $\sigma_{Ch} = 10^{-8} \text{ cm}^2$

- Limitation for d and γ :

$$kd < \gamma^2 (-0.125 \ln(\sigma_{Ch}/\sigma_T N^2))^{1/2}$$

- Attosecond bunches with width

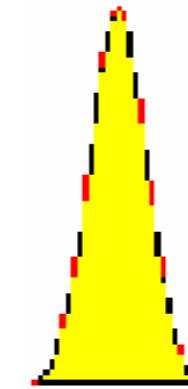
$$d \sim 1/k\gamma^2 \sim (100 \text{ as}) \cdot c$$

$$\gamma_{photon} = \frac{l}{4g^2} \text{ for } g = 100$$

$$g_{photon} = 40 \text{ keV}$$

$$\text{For } g = 10^3, g_{photon} = 6 \text{ MeV}$$

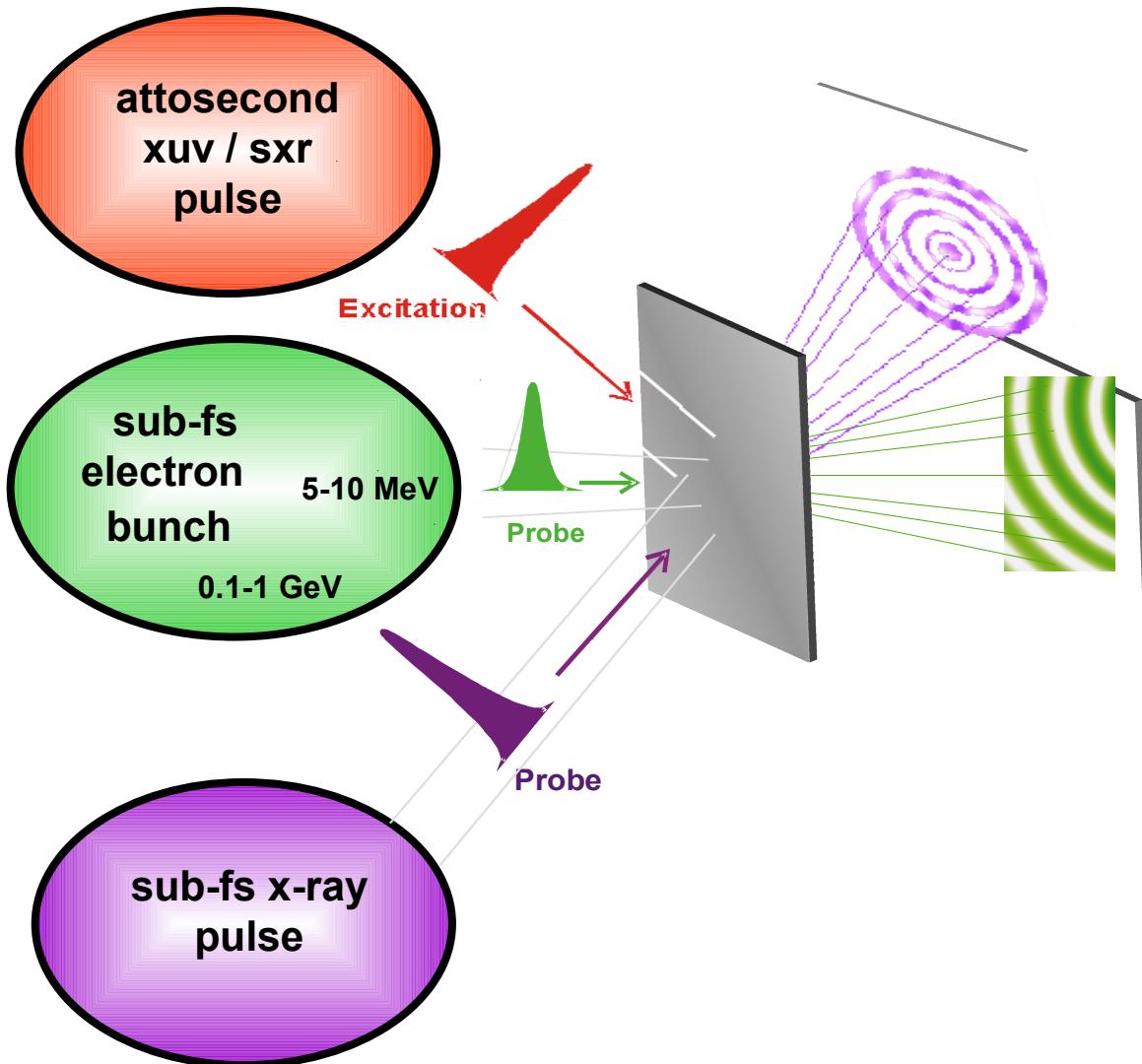
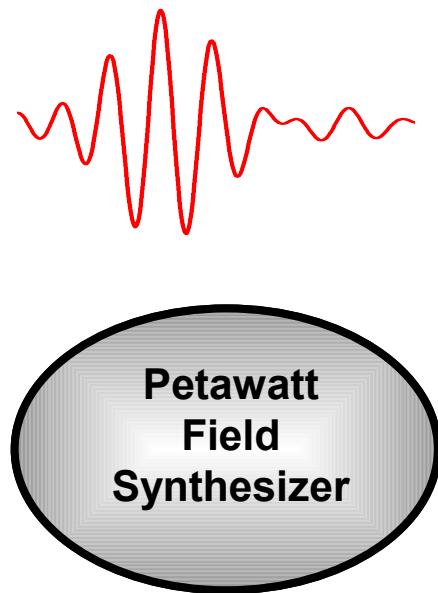
$\eta \sim 1$ efficiency



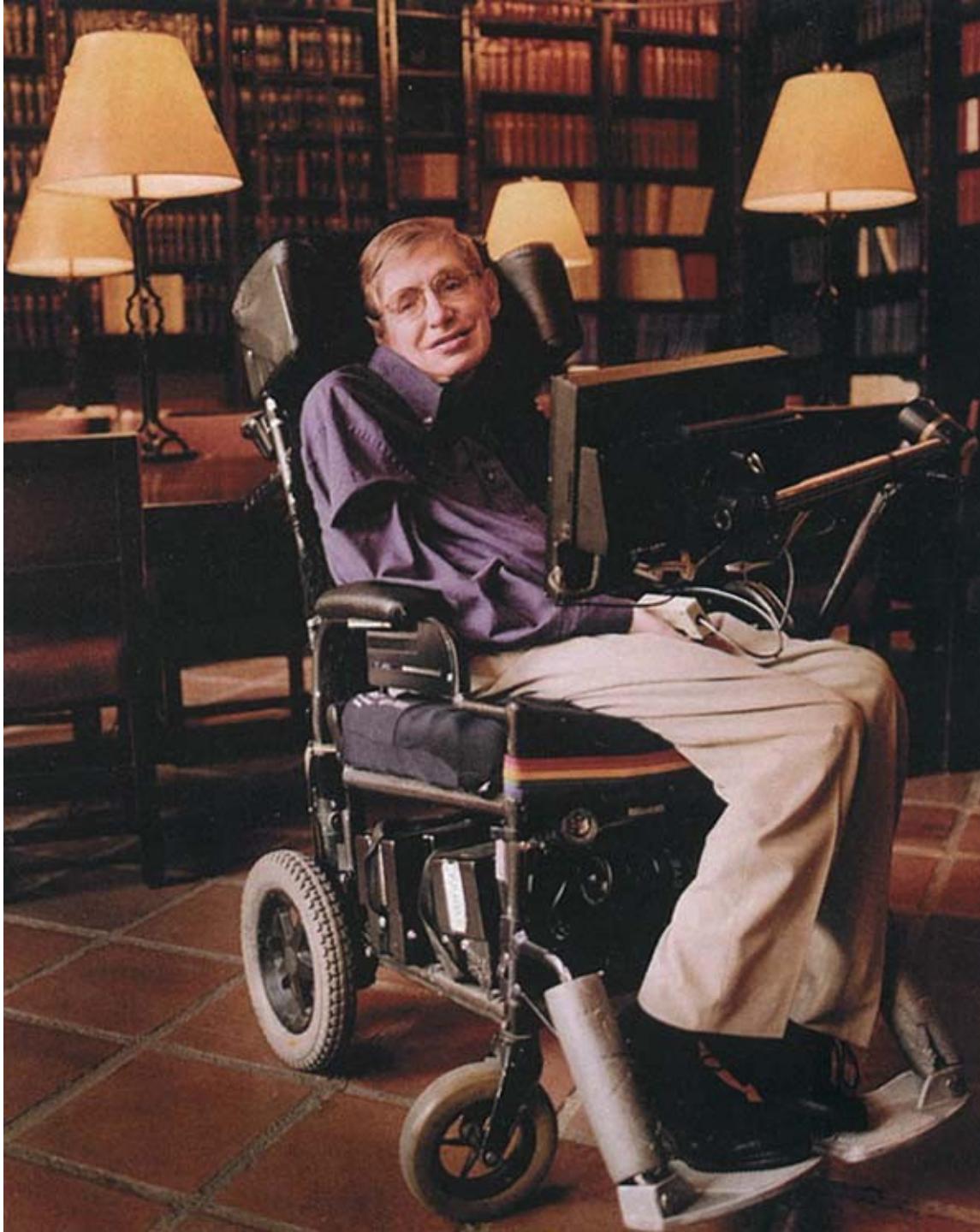
Bunch:
N particles
with
Gaussian
distribution

N. Naumova, I. Sokolov, J. Nees, A. Maksimchuk,
V. Yanovsky, and G. Mourou, Attosecond Electron Bunches,
Phys. Rev. Lett. **93**, 195003 (2004).

Control & 4D imaging of valence & core electrons with sub-atomic resolution

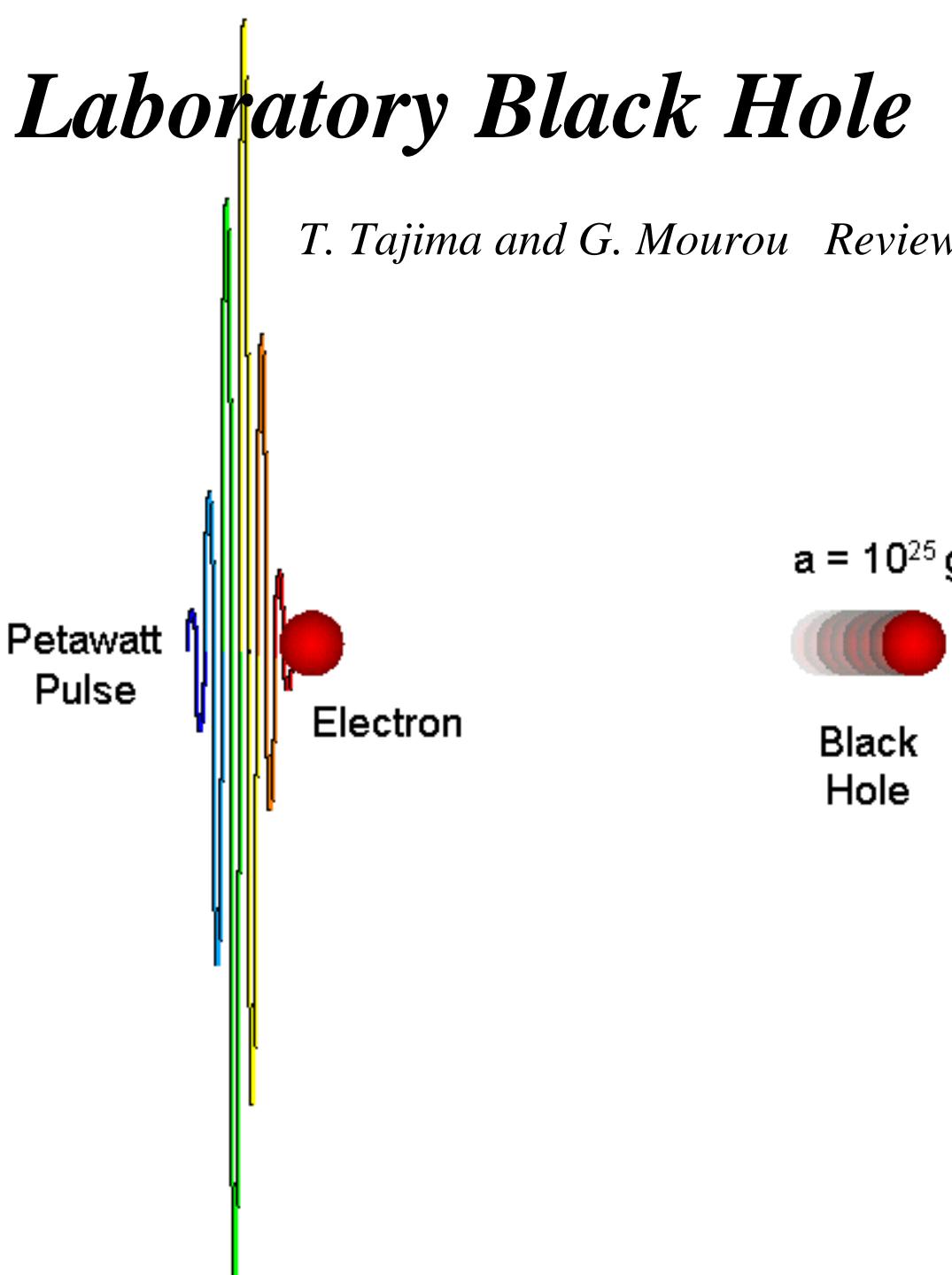


*Ultra-high Intensity
General Relativity
and Black Holes*



Laboratory Black Hole

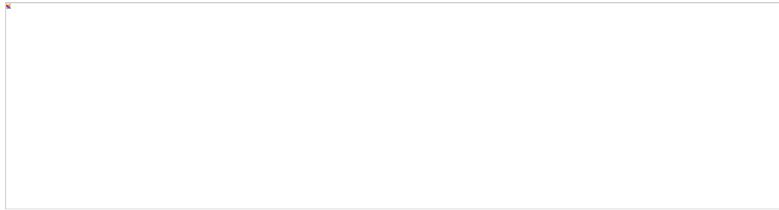
T. Tajima and G. Mourou Review of Modern Physics



*Equivalent to be near
a Black Hole of
Dimension?
Temperature?*

Is Optics in General Relativity?

Using the gravitational shift near a black hole:

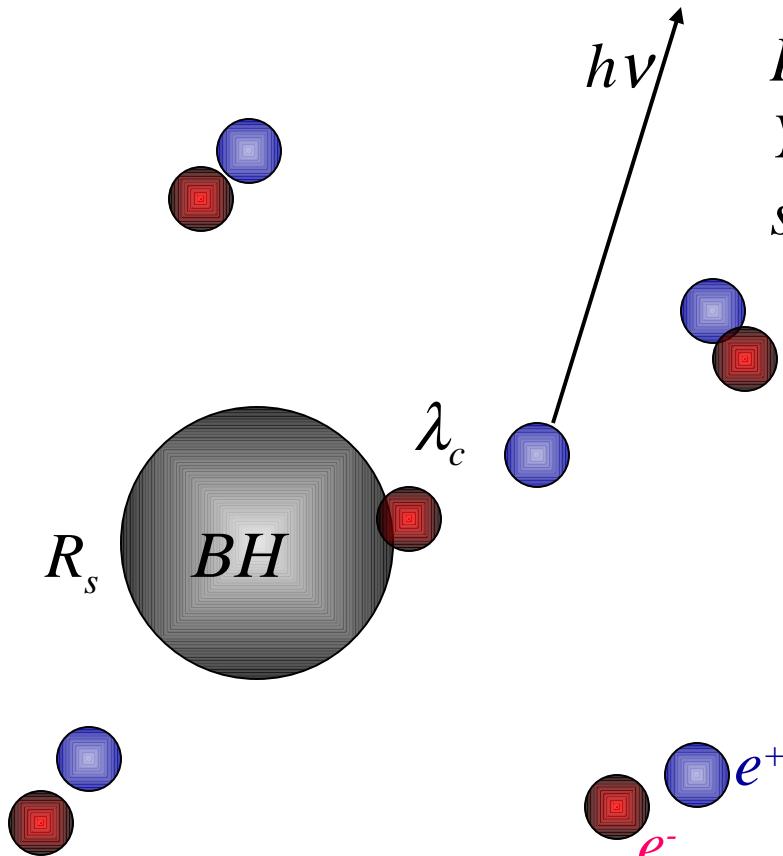


$$kT = \frac{ha_e}{2pc}$$

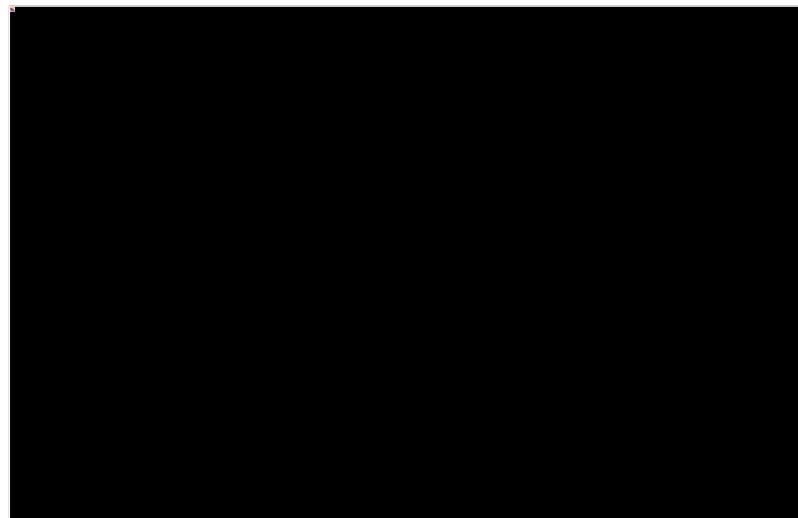


As we increase a_0 the Schwarzschild radius can become equal to the Compton wavelength.

Optics and General Relativity: Hawking Radiation



*In order to have Hawking radiation
You need the gravitational field
strong enough to break pairs*



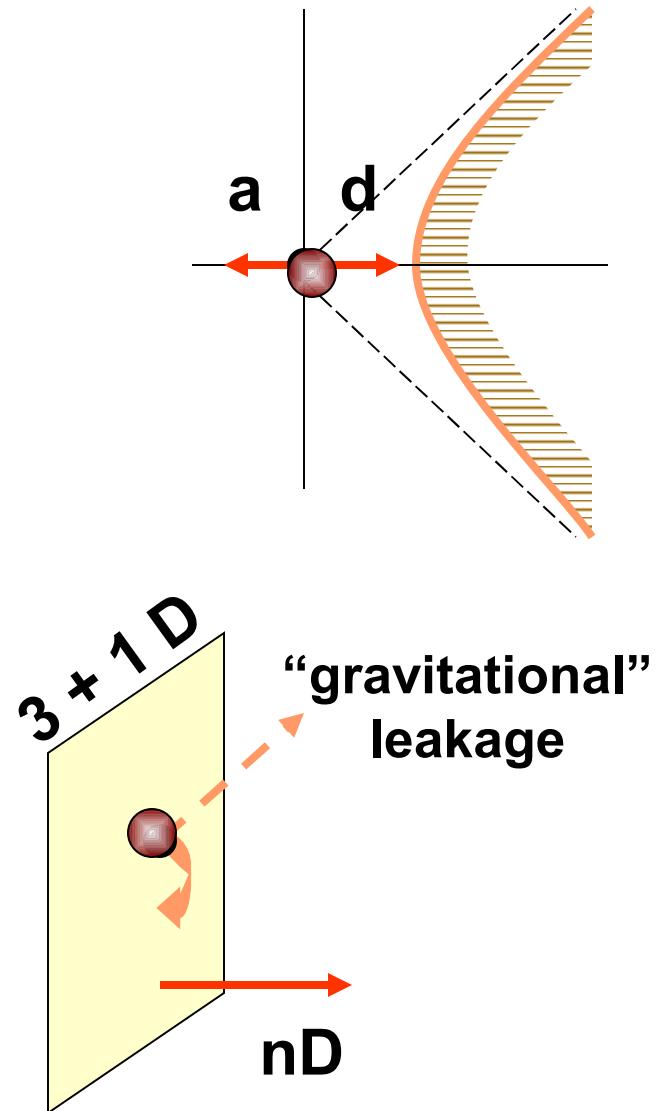
Finite Horizon and extra-dimensions

The distance to finite horizon is

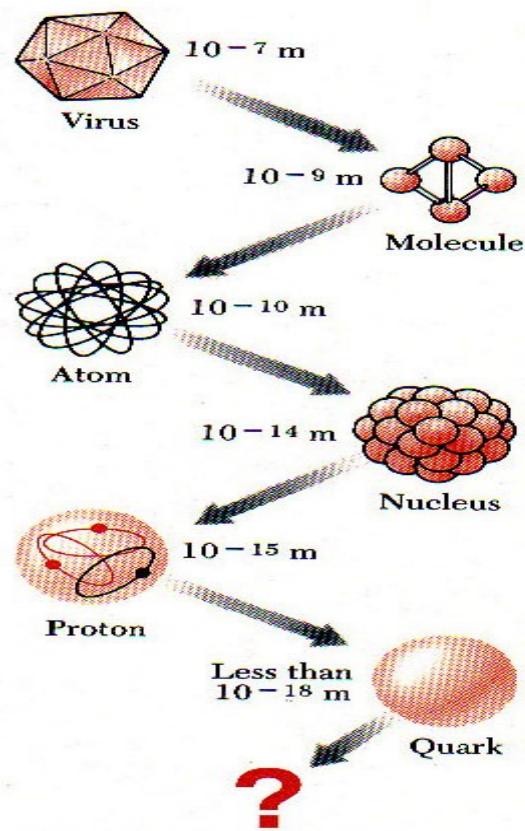


N. Arkani-Hamed et al. (1999)

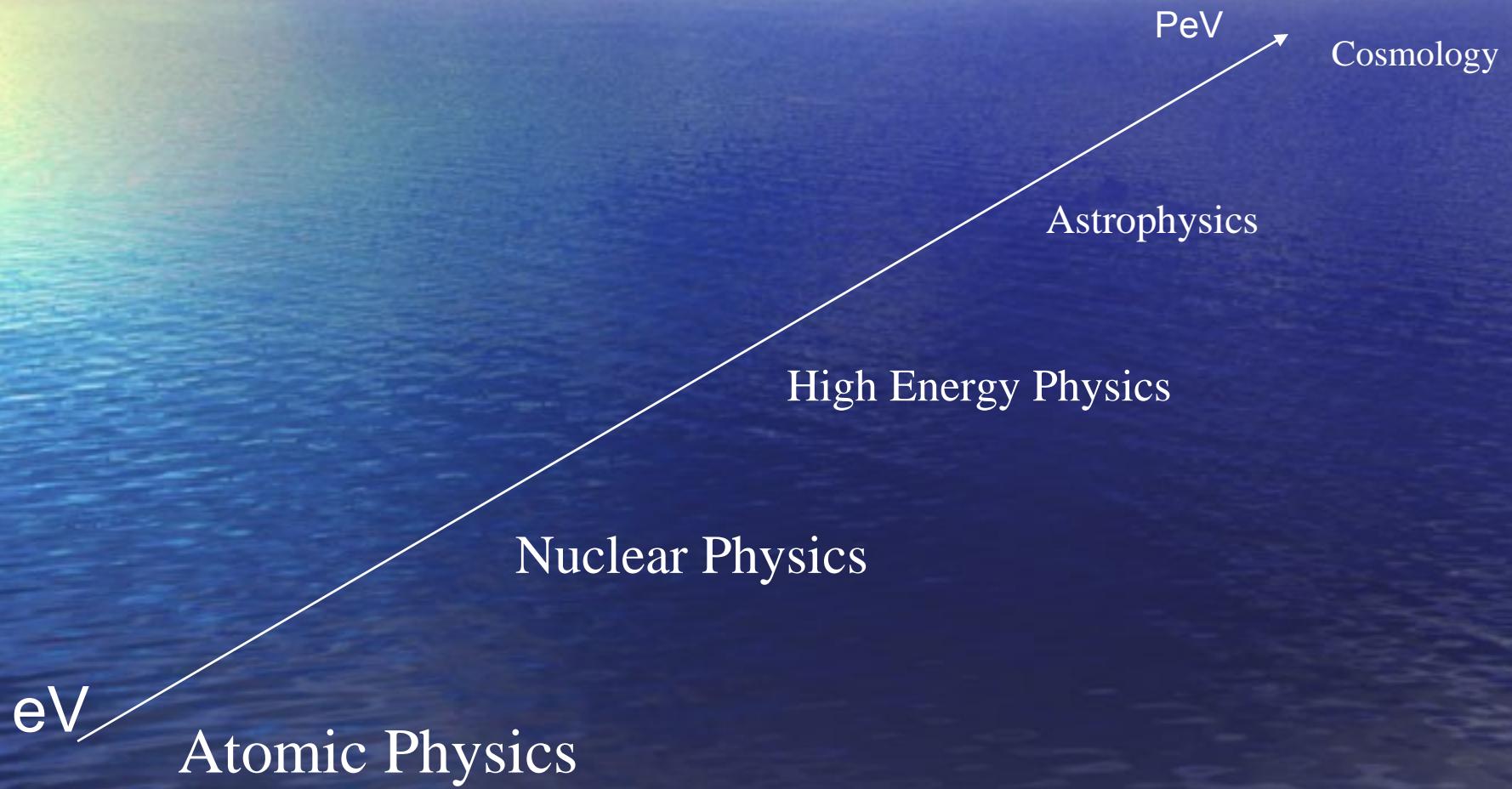
Up to $n=4$ extra-dimensions could be tested.



Moving from the Atomic Structure to the Quark Structure of Matter



Conclusion





Extreme Light Infrastructure

ELI

Extreme by its pulse duration

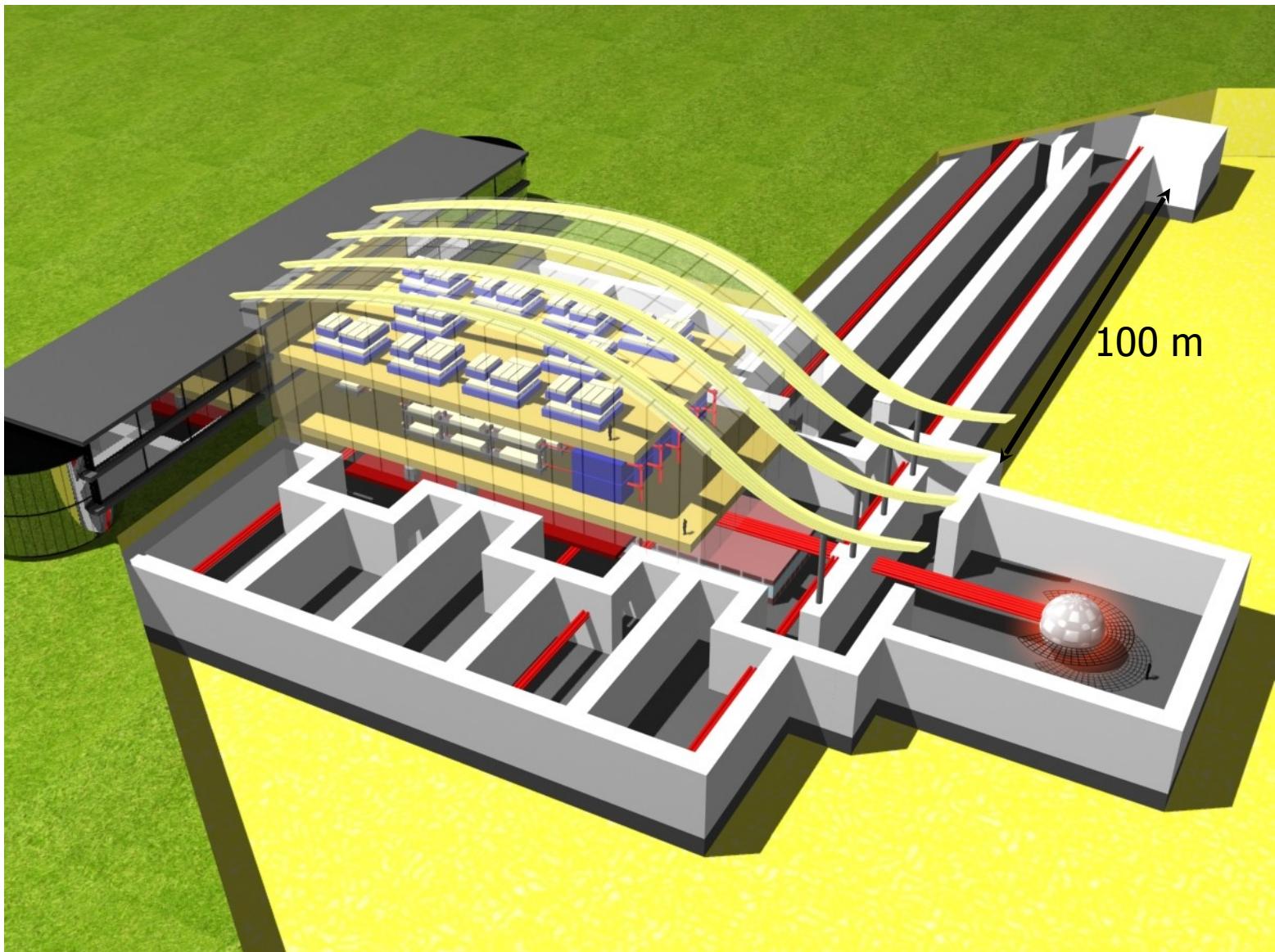
Extreme by its intensity

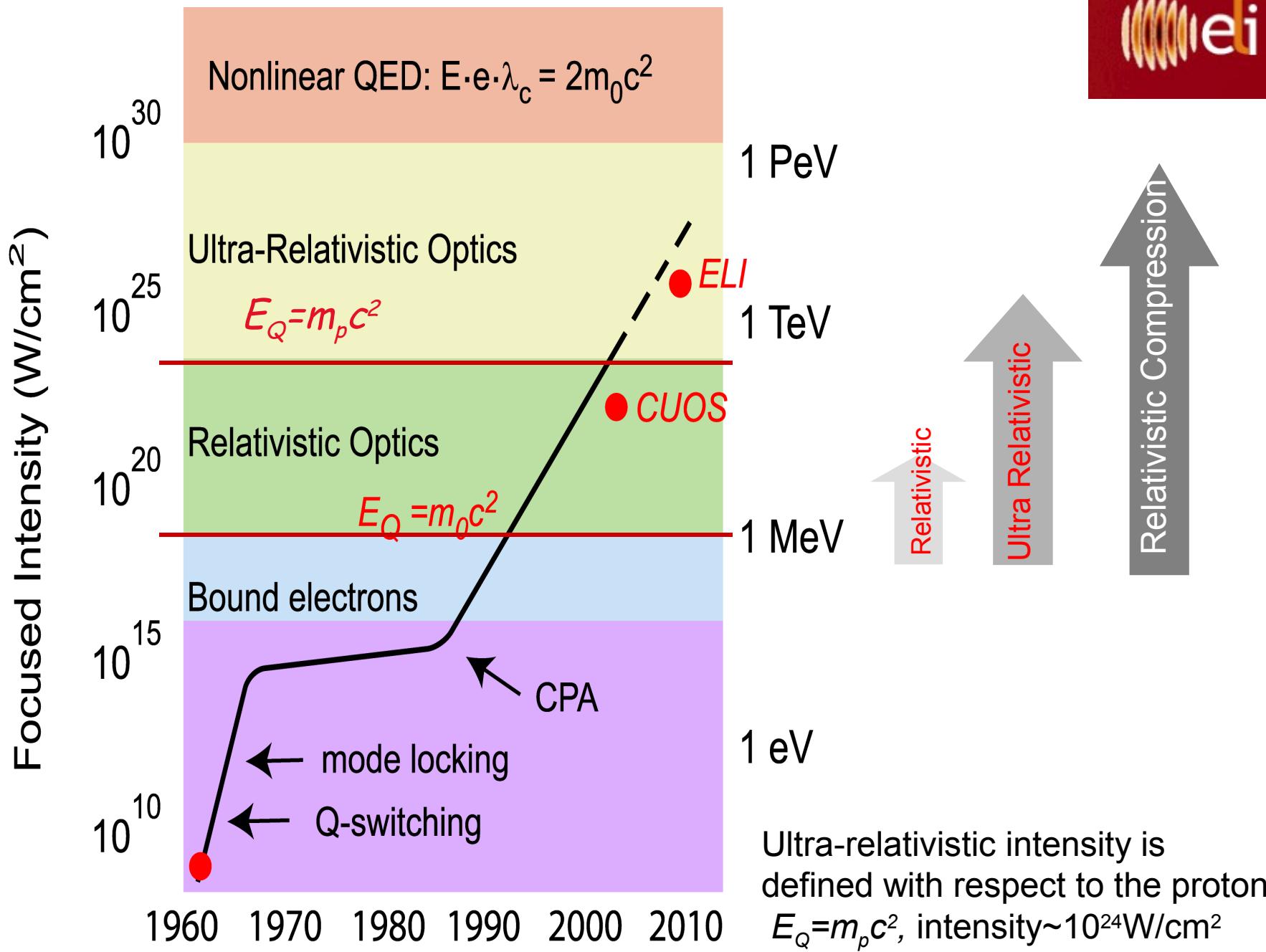
Extreme by the energy of its radiations and particles

On the Map of the Very Large
Scale European Infrastructures

ELI Proposal

<http://loa.ensta.fr>





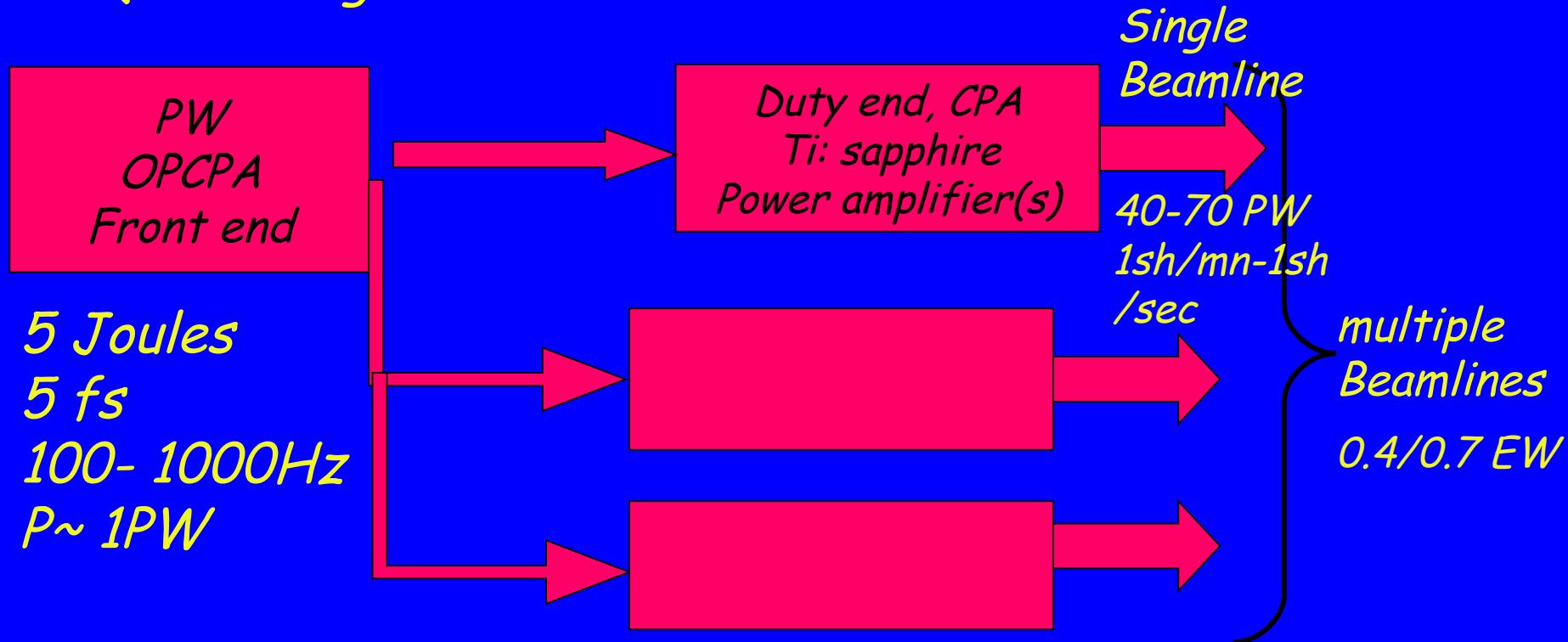
ELI Scientific Program: Based on an Exawatt Class Laser

Three Scientific Pillars

- **Ultra high Field Science:** access to the ultra-relativistic regime, ELI will afford new investigations in particle physics, nuclear physics, gravitational physics, nonlinear field theory, ultrahigh-pressure physics, astrophysics and cosmology.
- **Attosecond science:** snap-shots in the attosecond scale of the electron dynamics in atoms, molecules, plasmas and solids.
- **High Energy beam facility:** ELI will provide ultra-short energetic particle (>10 GeV) and radiation (up to few MeV) beams produced from compact laser plasma accelerators.

Exawatt laser scheme

MPQ Garching

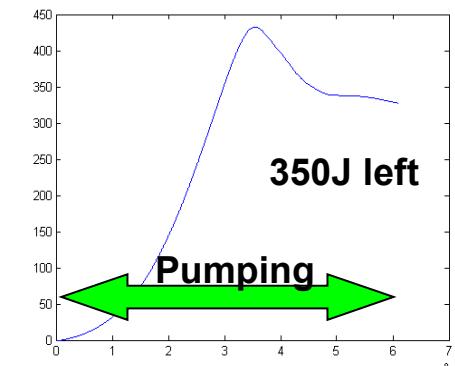
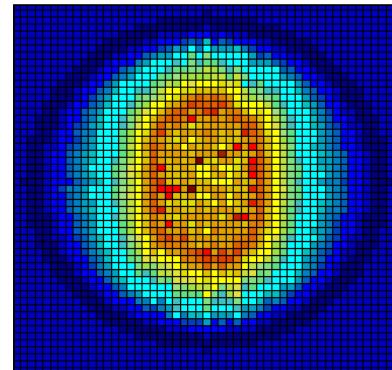


$$1 \text{ EW} = 1000 \text{ PW} = 10^{18} \text{ W} \rightarrow 10 \text{ KJ in } 10\text{fs}$$
$$0.1 \text{ EW} = 100 \text{ PW} = 10^{17} \text{ W} \rightarrow 1 \text{ KJ in } 10 \text{ fs}$$

Pumping of a Ø20cm crystal



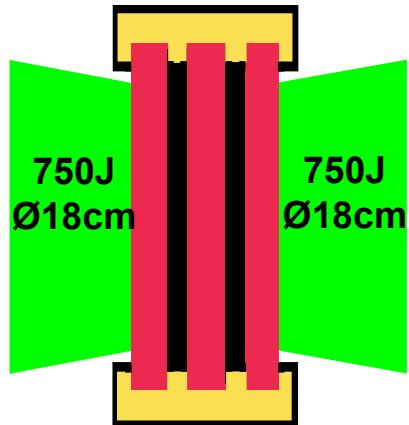
$$[\text{Ti}^{3+}] = 1.2 \cdot 10^{19} \text{ ions/cm}^3$$



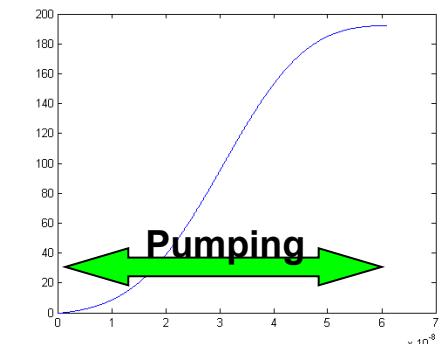
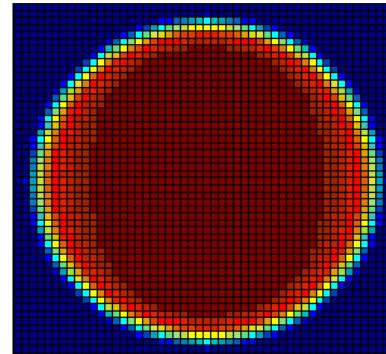
Gain profile after
pumping

Population inversion
versus time

Stack of low absorbing disks



$$[\text{Ti}^{3+}] = 6.6 \cdot 10^{18} \text{ ions/cm}^3$$



NO parasitic effects in crystals

Participating Countries

- France
- Germany
- United-Kindom
- Spain
- Italy
- Greece
- Lithuania
- Austria
- Romania
- Bulgaria
- Hungary
- Belgium
- Poland
- Portugal

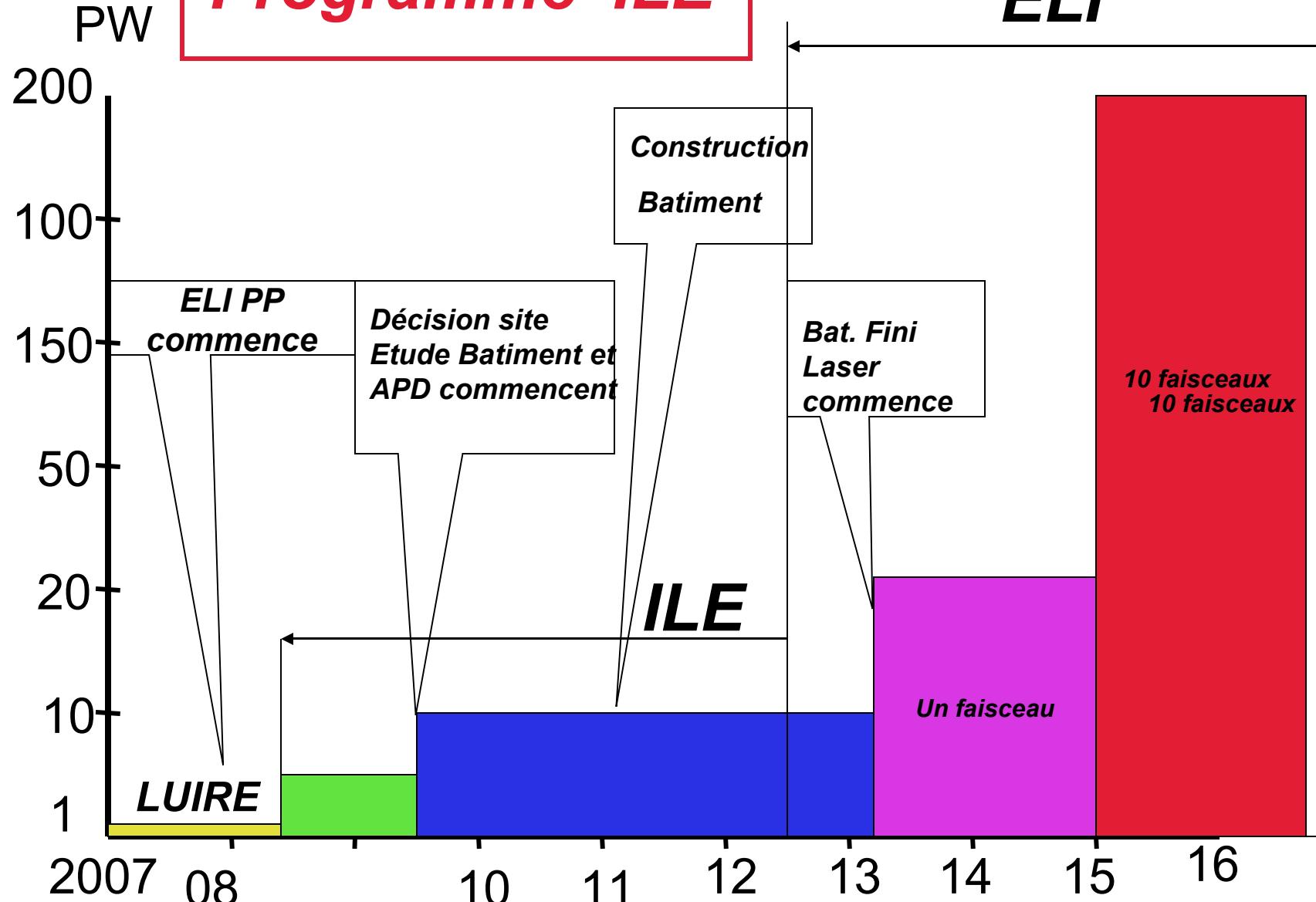


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- ²LOA, CNRS-ENSTA-X, 91761 Palaiseau (France)
- ³Johannes-Gutenberg-Universitat Mainz, D 55099 Mainz (Germany)
- ⁴GSI, Planckstr.1 64291 Darmstadt (Germany)
- ⁵LLR, CNRS-IN2P3-X, 91128 Palaiseau (France)
- ⁶Thalès, STI, 92704 Colombes cedex (France)
- ⁸General Physics Institute RAS, Moscow 119991 (Russia)
- ⁹IESL (Heraklion) 71110, Crete (Greece)
- ¹⁰Institut für Angewandte PM, D-85577 Neubiberg (Germany)
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- ¹²CPO, Bat. 101, 91898 Orsay cedex (France)
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- ¹⁵CENBG, CNRS, IN2P3, Université Bordeaux 1 (France)
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Programme ILE

ELI



Salle Jaune épaisseur du
trait

Thank you

Become an ELI enthusiast
You can register @

WWW.eli-laser.eu
Eli@eli-laser.eu

*One of the big Challenges
in Physics would be to built
A laser powerful enough to
breakdown vacuum.*

Survey by Science 2005



Towards the Exawatt Laser: Relativistic and Ultra-Relativistic Optics

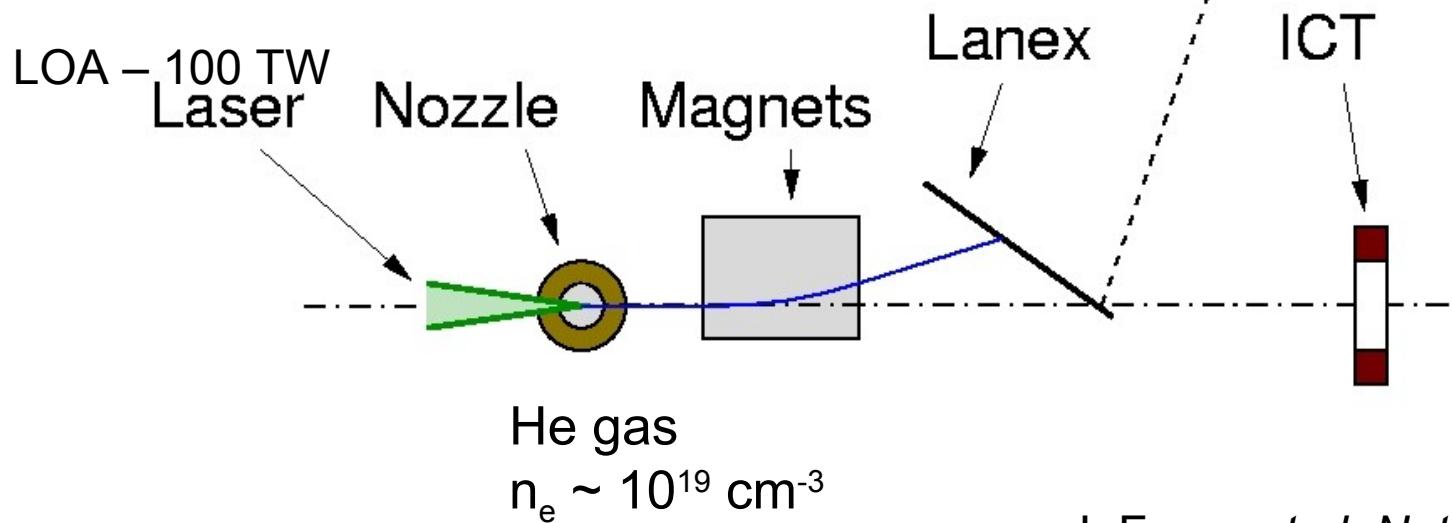
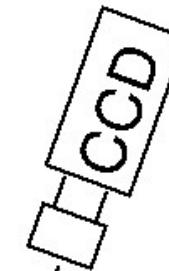
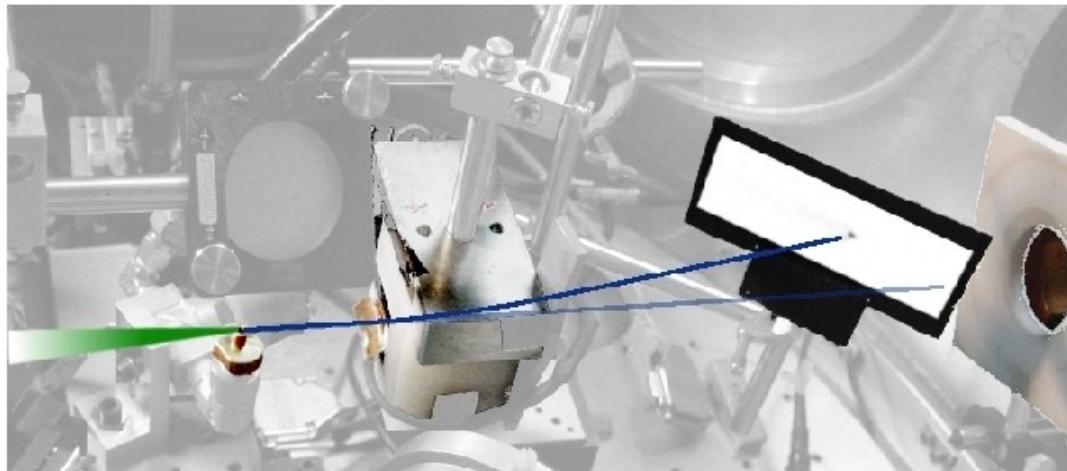
Gérard A. MOUROU

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gerard.mourou@ensta.fr

CLEO Europe 2007

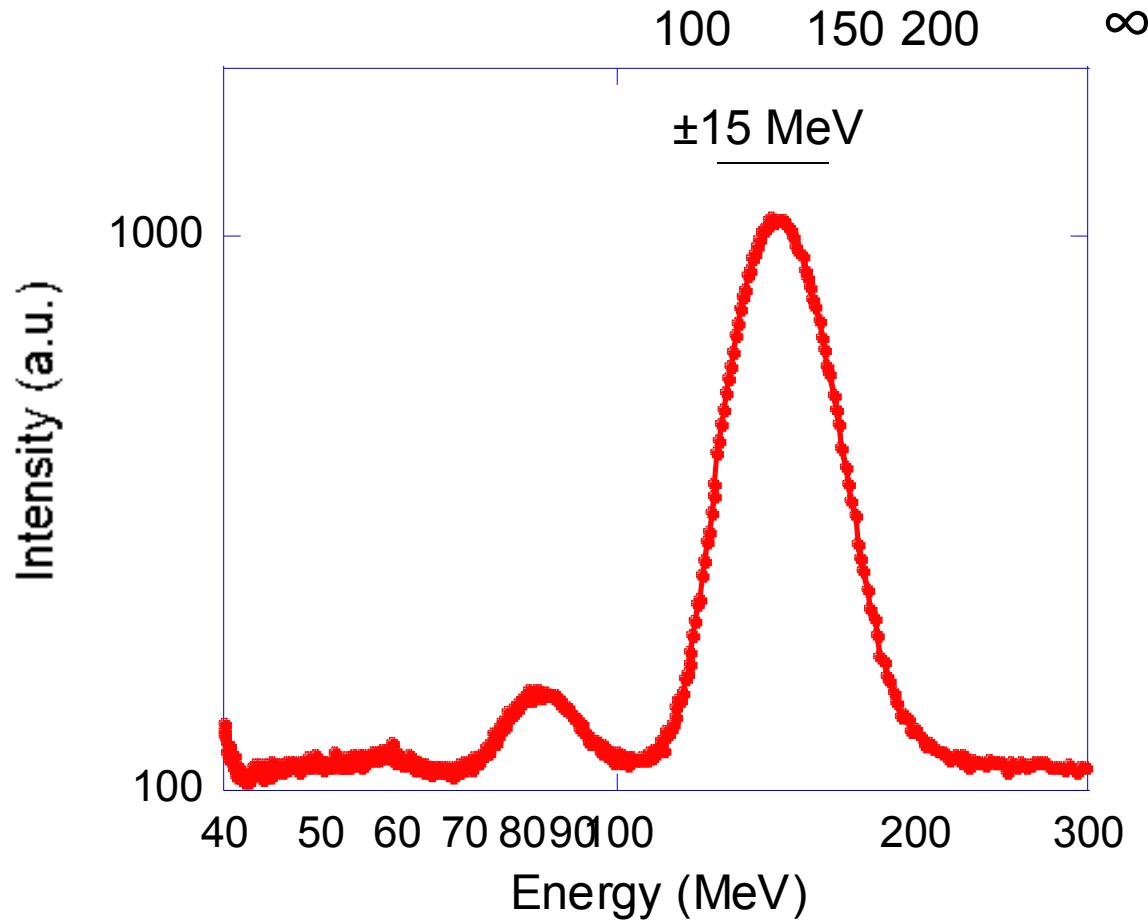
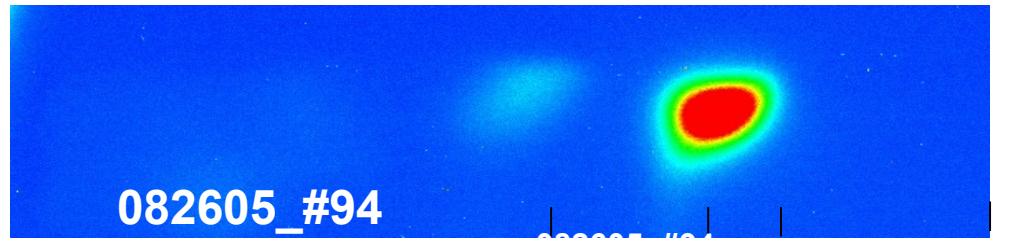
Recent results on electrons acceleration - Setup



J. Faure *et al*, *Nature* 2004



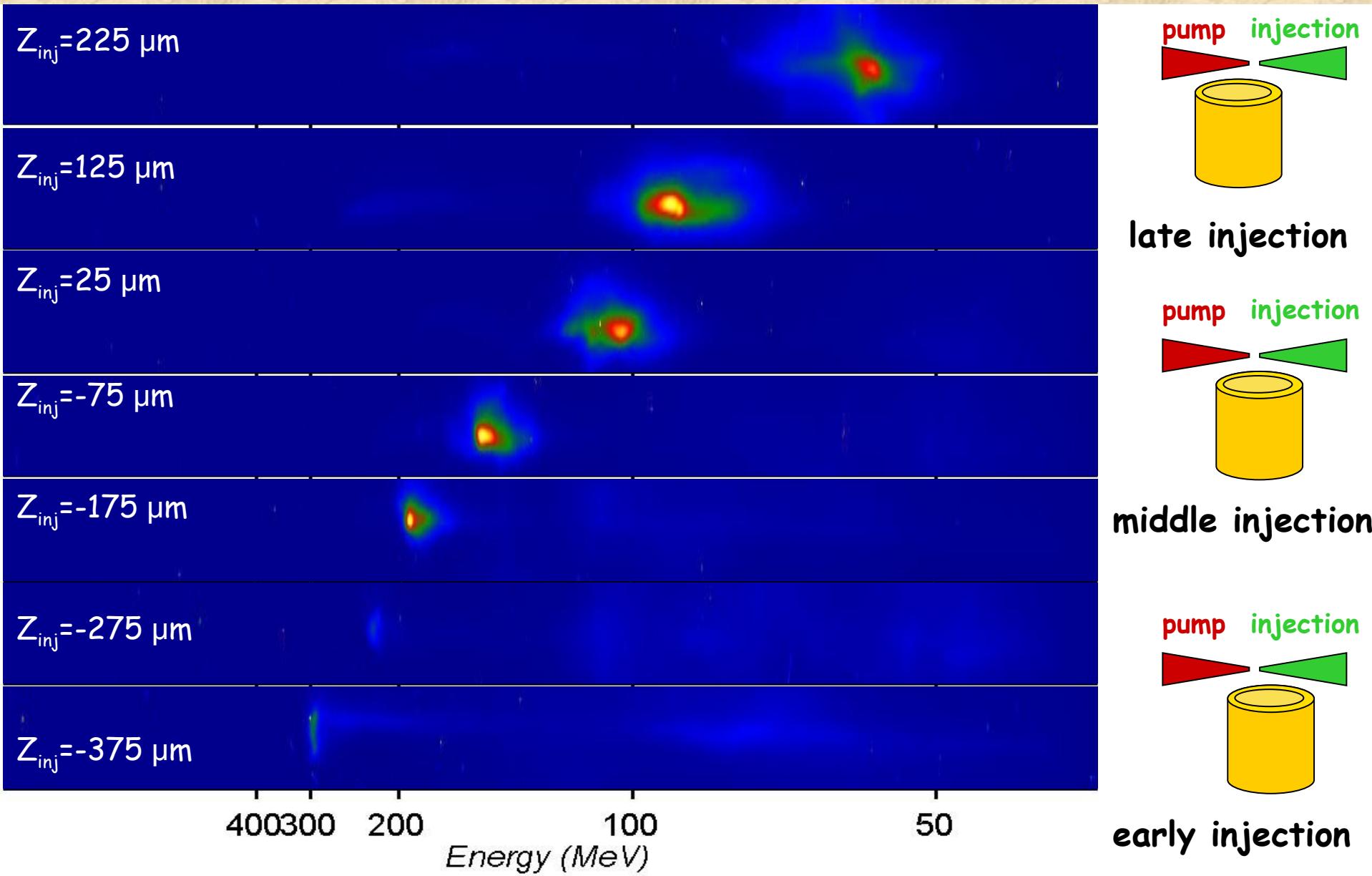
Quasi-monochromatic beam with $E_{\max} = 160 \text{ MeV}$ at $n_e = 2.10^{19} \text{ cm}^{-3}$





Tunable monoenergetic bunches

V. Malka and J. Faure



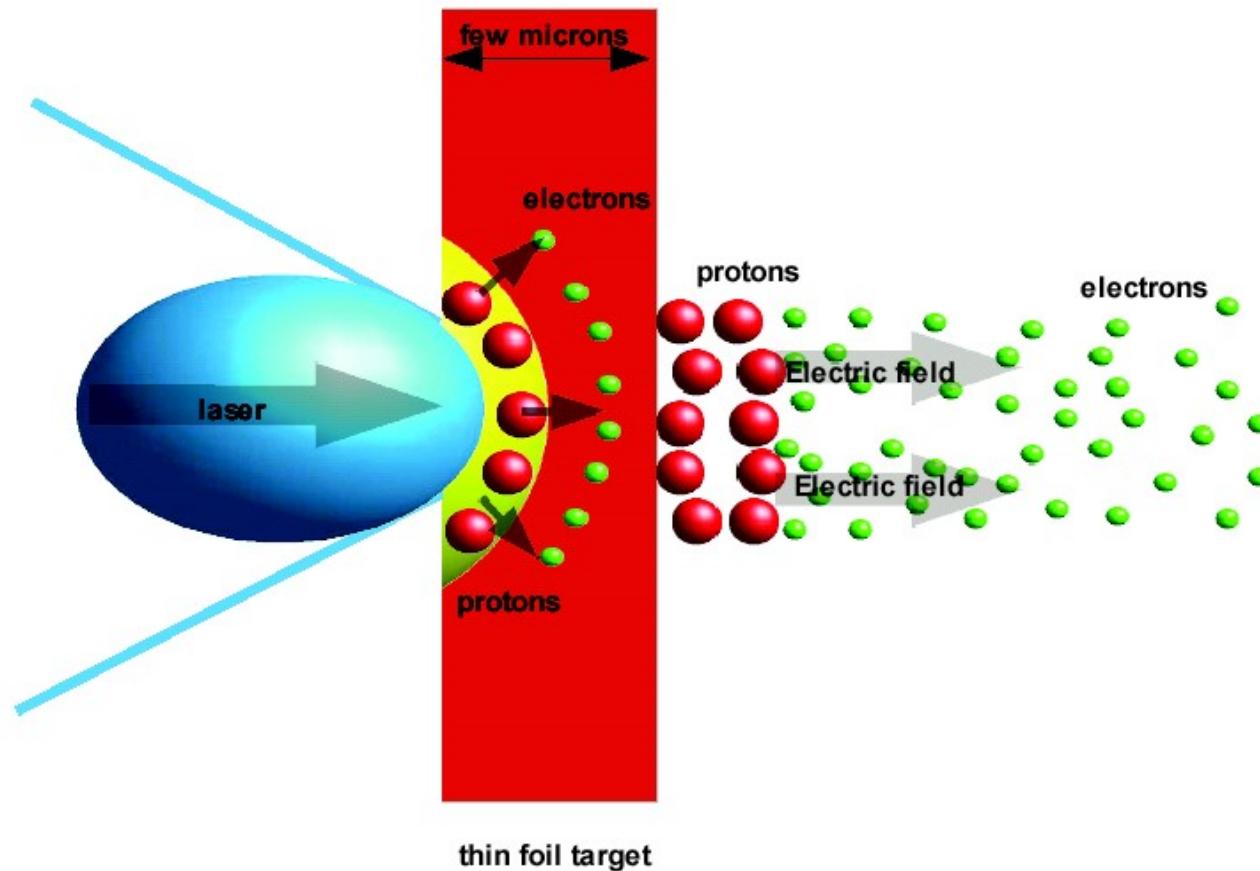
QuickTime^a and a
Photo - JPEG decompressor
are needed to see this picture.

T. Tajima and J. M. Dawson,
Laser electron accelerator,
Phys. Rev. Lett. 43, 267 (1979)

Secondary effects of
electron acceleration:

Proton Acceleration

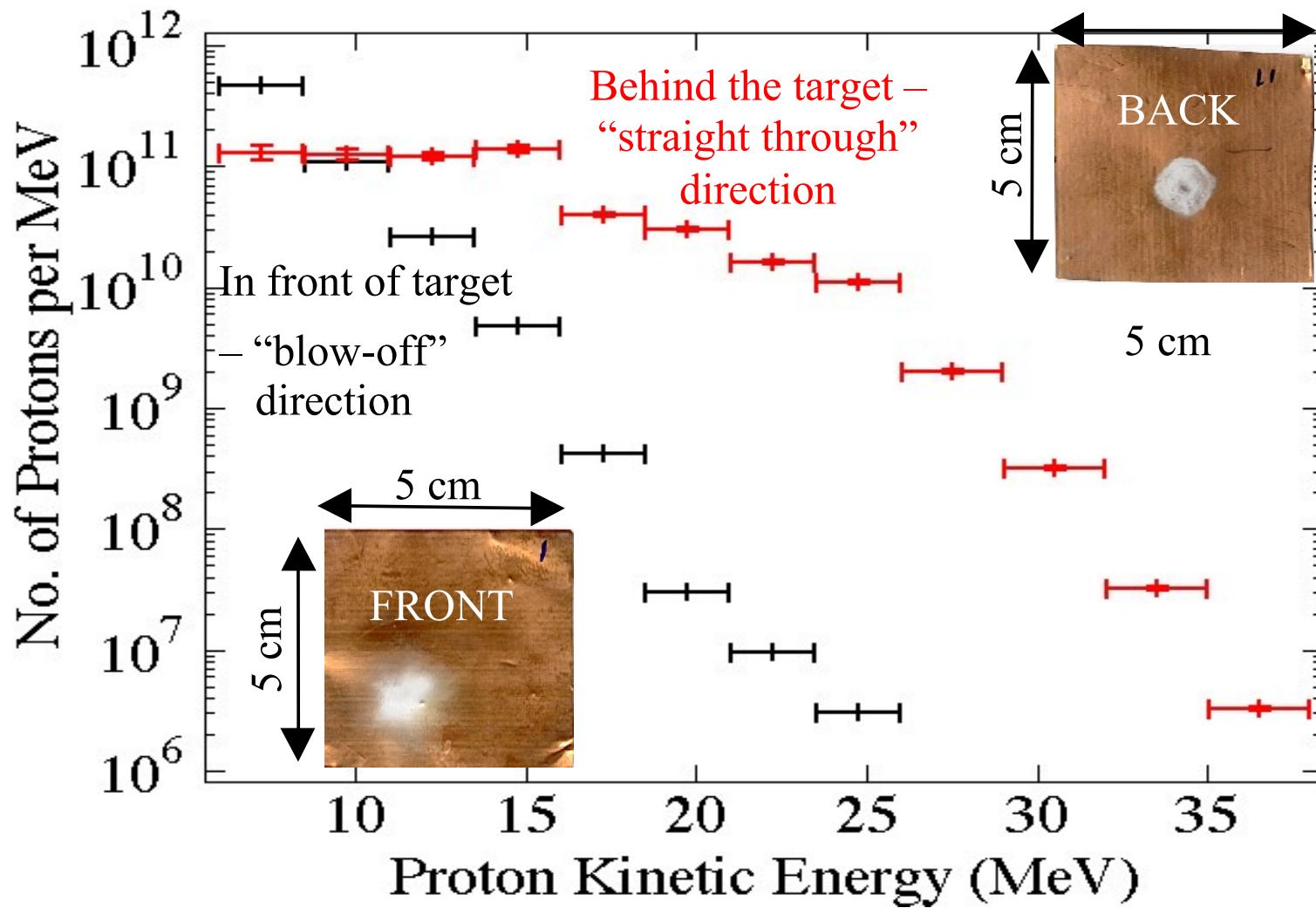
Front and back acceleration mechanisms



Peak energy scales as : $E_M \sim (I_L \times \lambda)^{1/2}$

Large Laser results : Vulcan laser

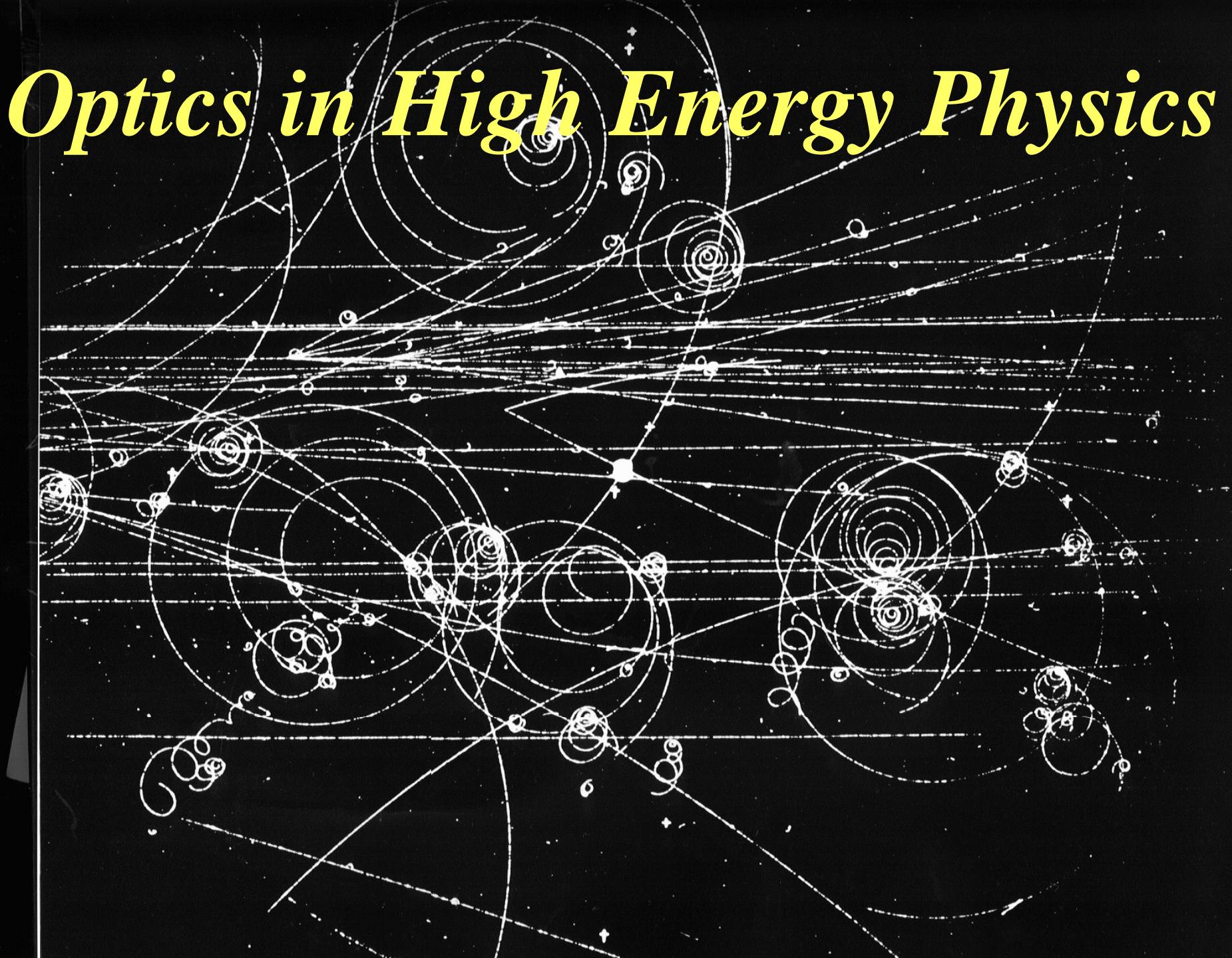
50J:1ps & 1shot/20min.



Secondary effects of
electron acceleration:

X-rayBeam

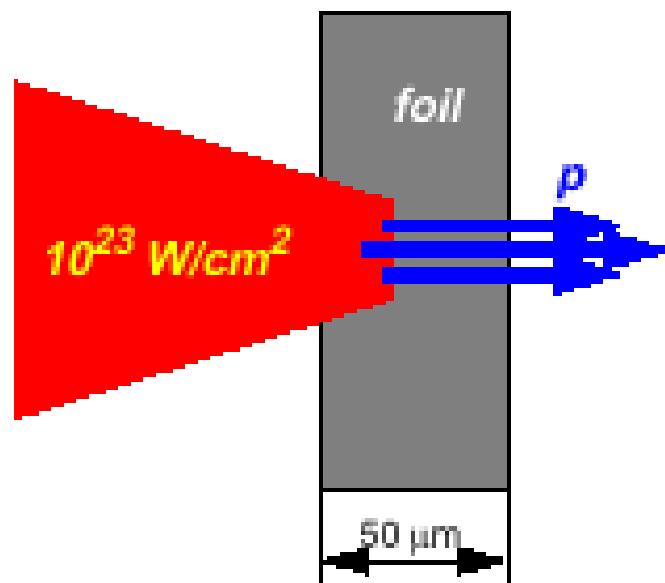
Optics in High Energy Physics



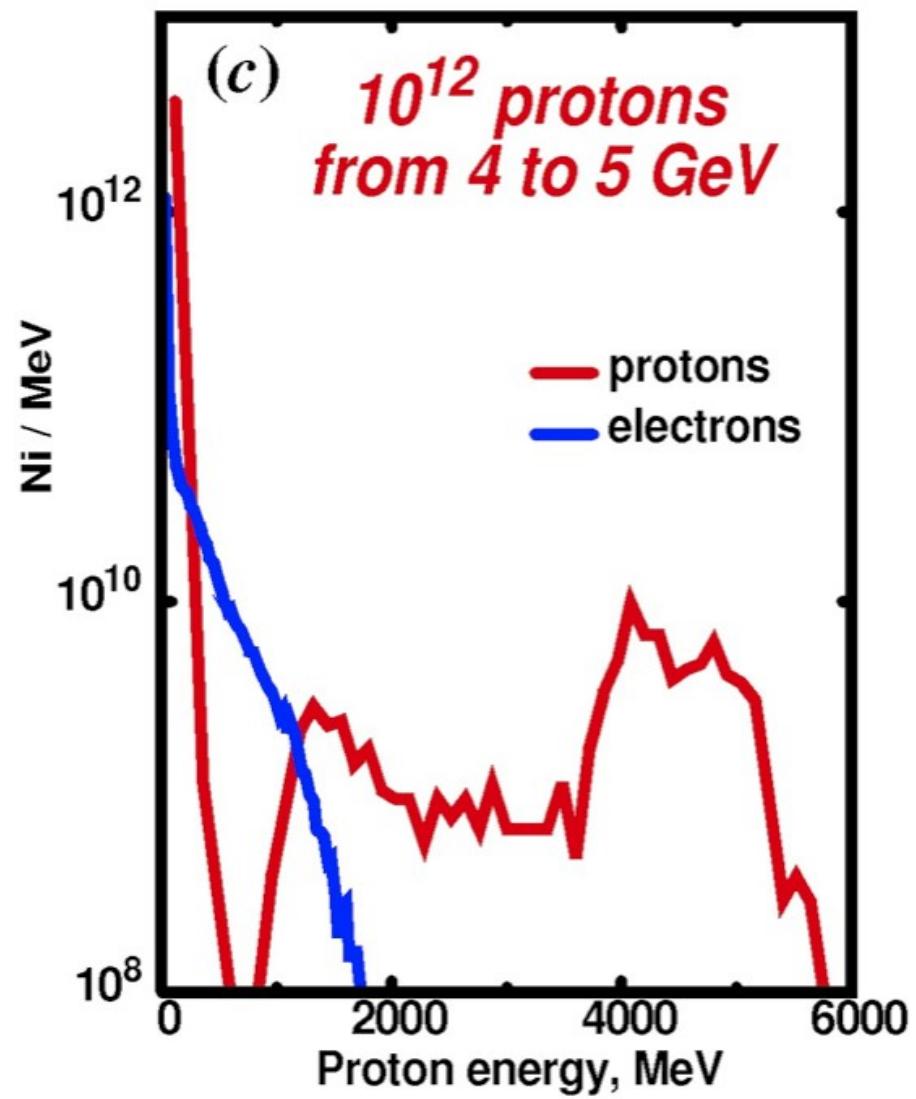
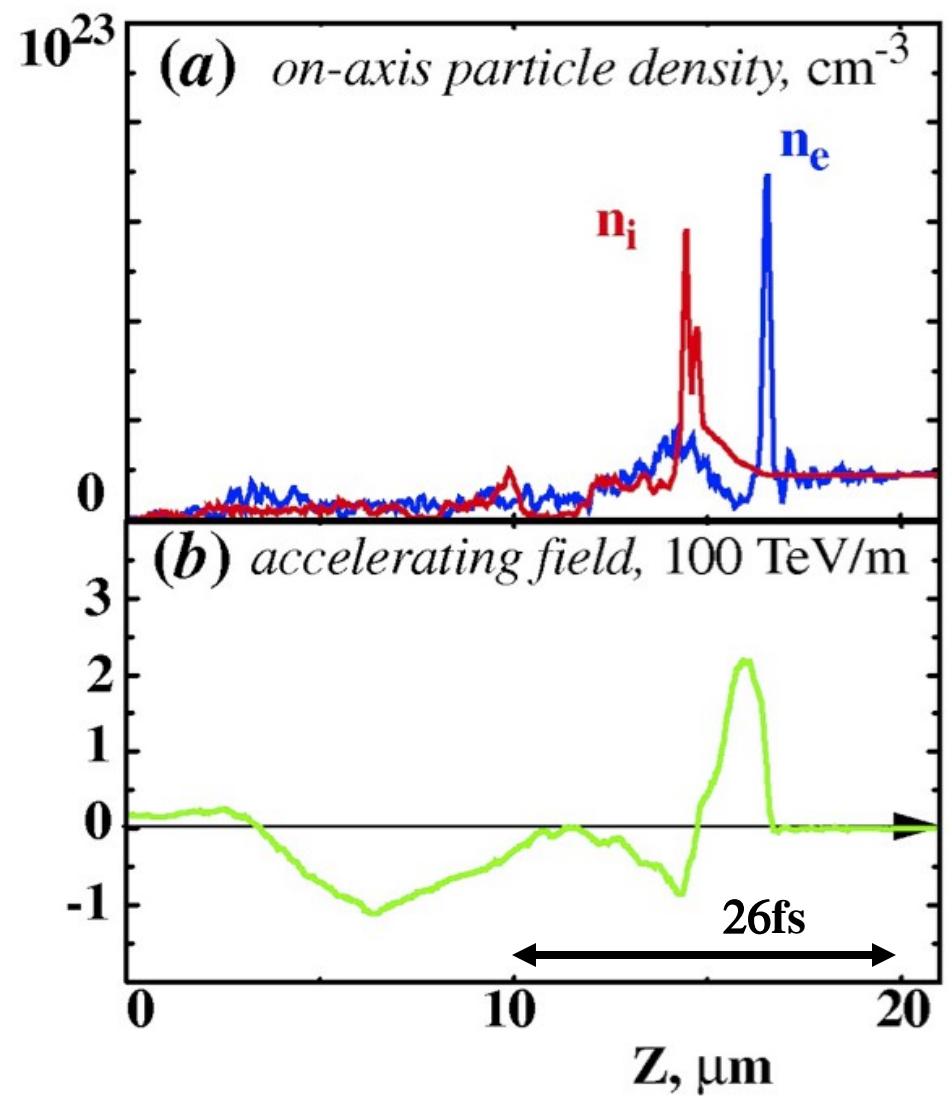


5 GeV proton bunch at solid state density

3d PIC simulations, A.Pukhov, Theorie, MPQ,

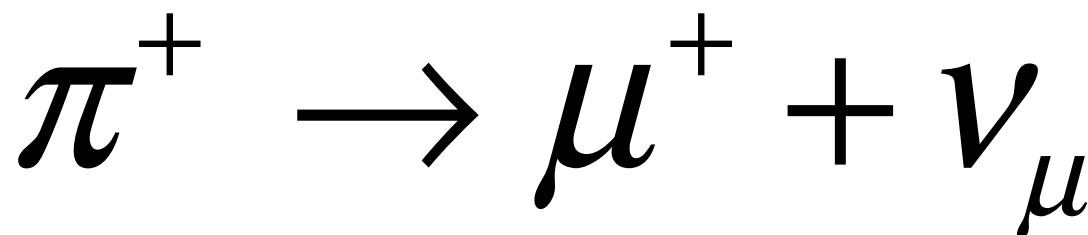


1 kJ, 15 fs long laser pulse
focused down to 10 μm spot
on a plasma with $n=10^{22} \text{ cc}^{-1}$



Unstable Particle Acceleration Muon and neutrino Beams

G. Mourou, S. Bulanov, T. Tajima Review of
Modern Physics 2006



Pions have 20ns lifetime (6m). They can only be accelerated up to 100MeV during this time with conventional technology. Their mass is $\sim 200\text{MeV}$, to increase their lifetime 100times, to $2\mu\text{s}$, we need to increase their energy by 100 to 20GeV. This can be achieved with laser technology over only $20\mu\text{m}$.